


# Phase transitions in dilute stellar matter

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 Supernova remnant  
and *neutron star* in  
Puppis A  
(ROSAT x-ray)

**Motivation:** study of  
dishomogeneous phases in PNS  
crusts and supernova cores.

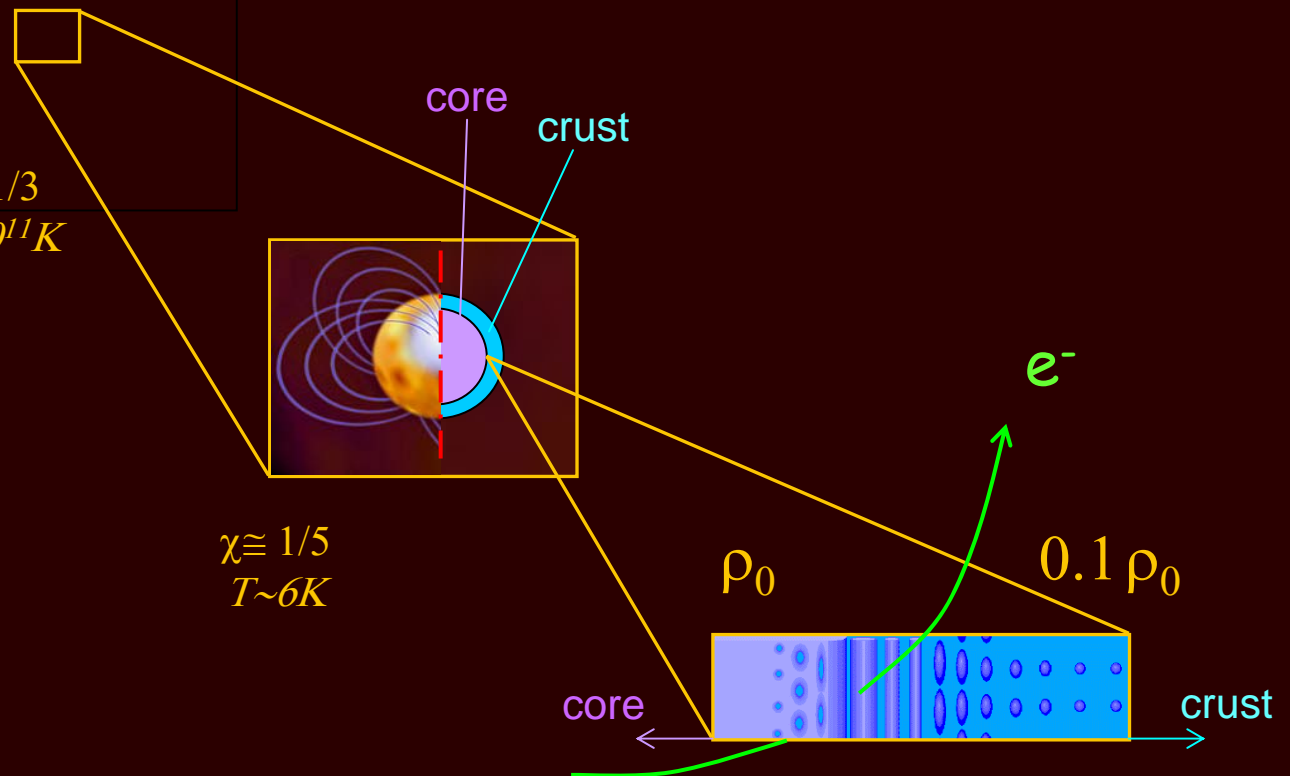
Influence on the explosion dynamics  
and the PNS cooling via:

- **electron capture rates** *Janka et al PR442 (2007)*
- **neutrino opacity** *H.Sonoda PRC75(2007)*
- **heat capacity** *D.Page NPA777(2006)*

$\chi \cong 1/2$   
 $T \sim 10^{12} K$

$\chi \cong 1/3$   
 $T \sim 10^{11} K$

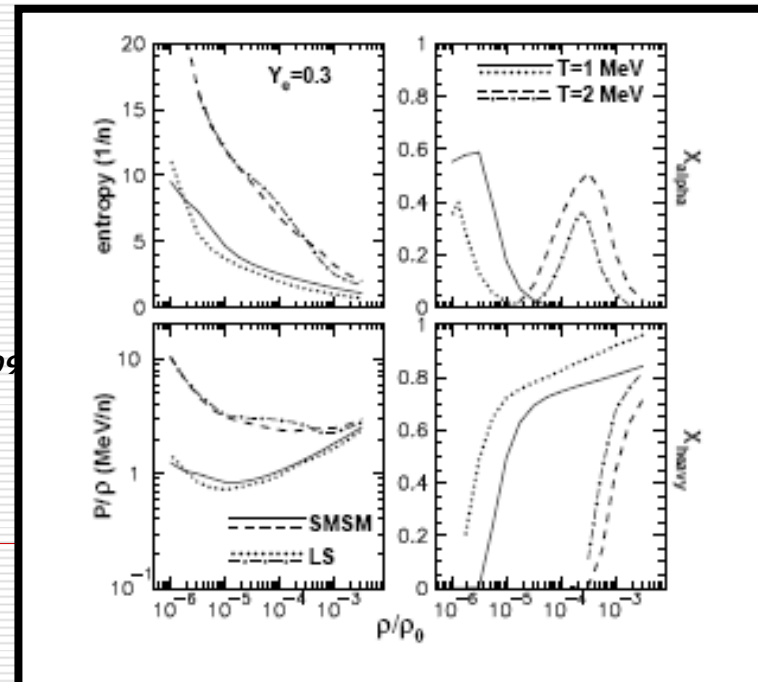
$\chi \cong 1/5$   
 $T \sim 6K$



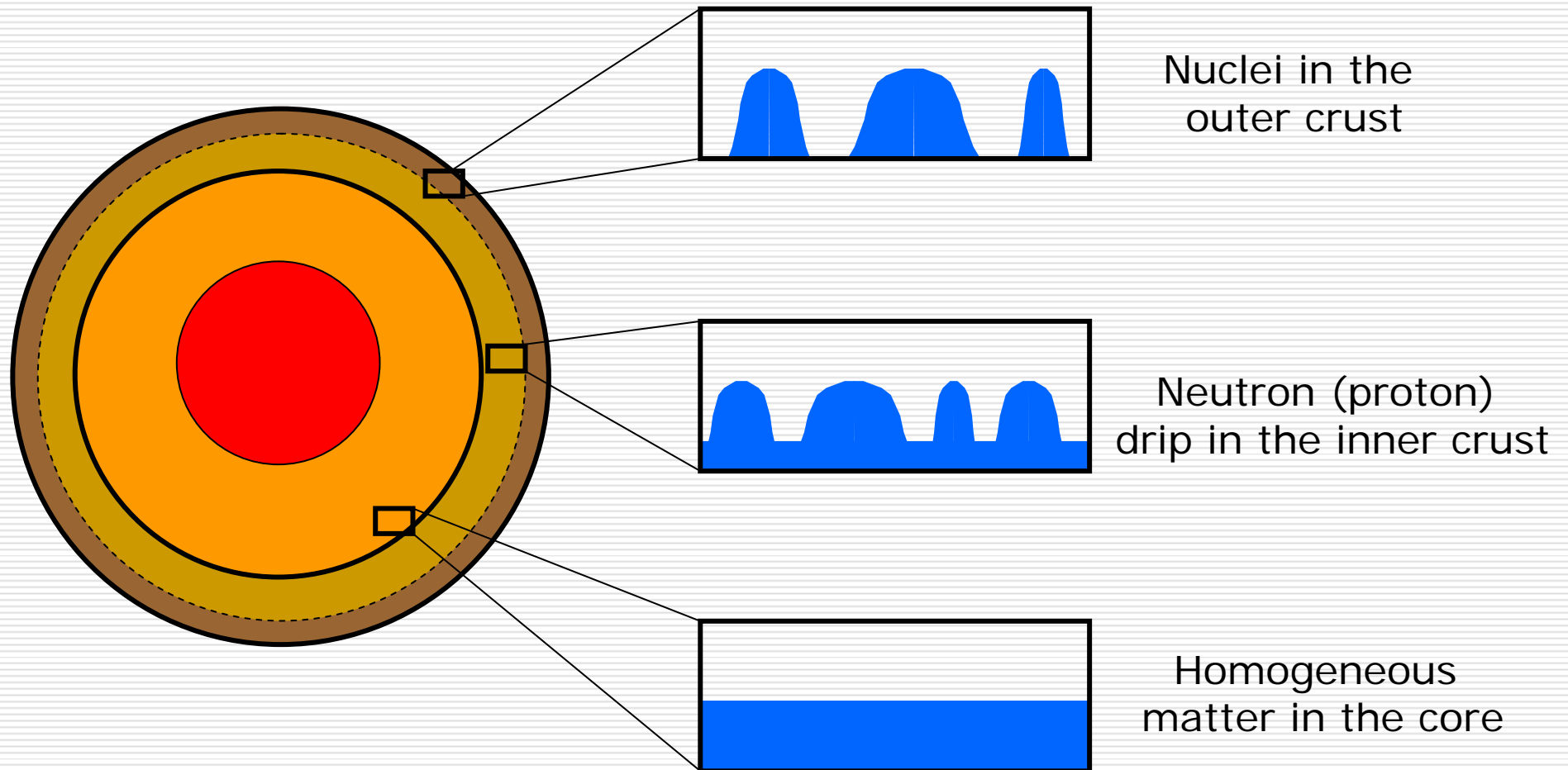
# Dilute stellar matter at $T > 0$

- Standard treatment in supernova codes: statistical equilibrium of  $n, p, \alpha, 1$  heavy cluster + 1st order phase transition to uniform matter  
*Lattimer-Swesty EOS, Shen EOS*
- **Improvement:** non-interacting ideal-gas of nuclei (NSE) *R.S.Souza et al Astrophys.J.707:1495-1505,2009 A.Botvina et al Nucl. Phys. A:98-132,2010 S.I.Blinnikov et al 0904.3849*
- **Missing physics:** in medium corrections, inter-particle interactions

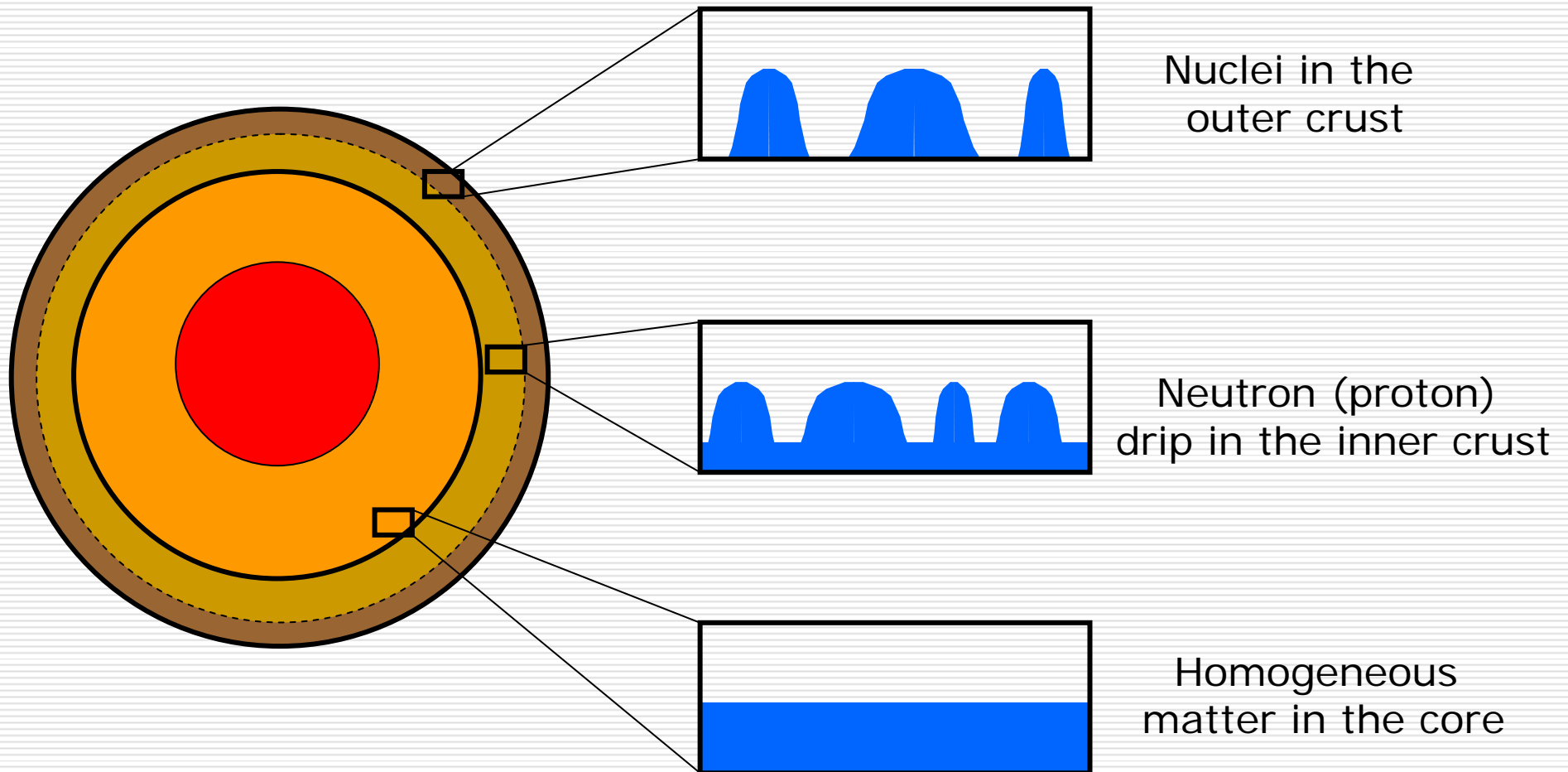
- Interactions in the S matrix formalism  
*S.K.Samaddar et al. Phys.Rev.C80:035803,2009*
- Virial EOS  $A < 5$  *A.Schwenk et al. Phys.Rev.C78:015806,2008*
- Quasi-particle gas model  
*Skyrme +  $A < 14$  S.Heckel et al Phys.Rev.C80:015805,2009*  
*RMF +  $A < 4$  S.Tygel et al Phys.Rev.C81:015803,2010*
- Phenomenological models  
*M.Hempel et al., astro-ph/0911.4073*



# Dilute matter at $T > 0$ : the phenomenology of the crust-core transition

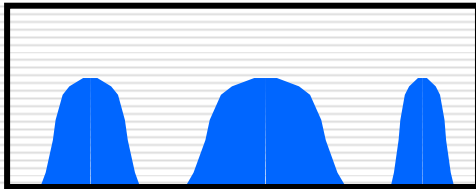


# Dilute matter at $T > 0$ : the phenomenology of the crust-core transition

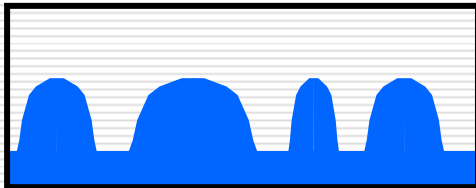


+ electrons  $\rho_e = \rho_p$ ,  $\gamma, \nu$

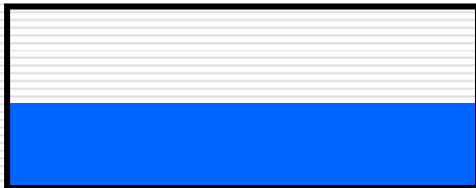
# An hybrid model for the crust-core transition



Nuclei in the outer crust



Neutron (proton) drip in the inner crust

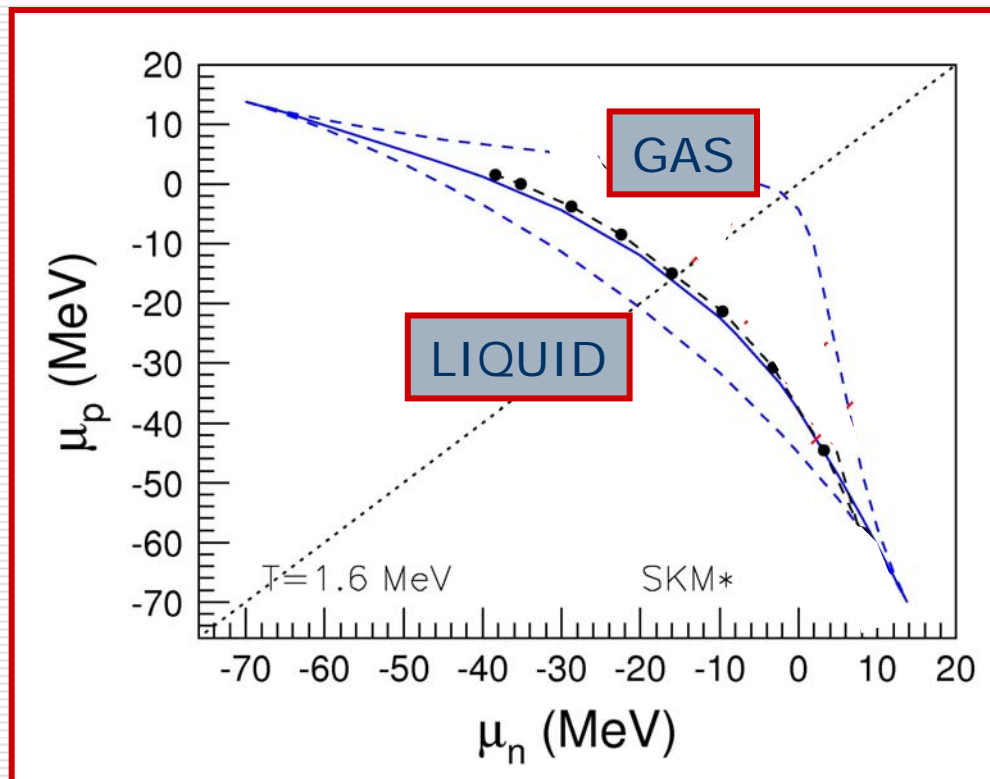


Homogeneous matter in the core  $\rightarrow$  Finite temperature Hartree F with Skyrme interactions (SI

$$\rho_{p,n} = \frac{2}{h^3} \int_0^{\infty} \frac{4\pi p^2 dp}{1 + e^{\beta(p^2/2m_{p,n}^* + U_{p,n} - \mu_{p,n})}}$$

+ electrons  $\rho_e = \rho_p, \gamma, \nu$

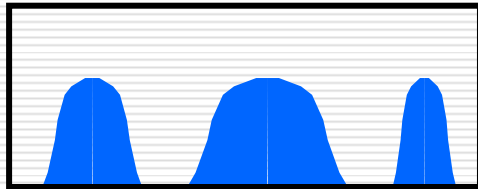
# Thermodynamics of homogeneous matter in the MF approximation



- Coexistence of purely homogeneous NM
- ⋯ Spinodal of purely homogeneous NM

=> Homogeneous matter is unstable over a wide  $\mu, T$  region

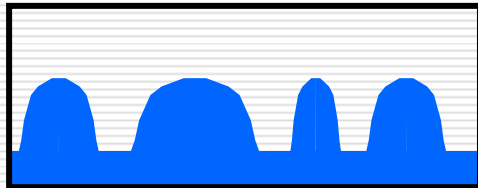
# An hybrid model for the crust-core transition



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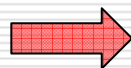
Statistical ensemble of  
interacting excited clusters



Neutron (proton)  
drip in the inner crust



Homogeneous  
matter in the core



Finite temperature Hartree Fock  
with Skyrme interactions (SKM\*, Sly230a)

+ electrons  $\rho_e = \rho_p$ ,  $\gamma, \nu$



# Statistical ensemble of interacting excited clusters

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## □ Standard NSE

$$F_{i,n_i} = -Tn_i (\ln Z_i + 1) \quad Z_i = \frac{g_i(T)V}{n_i} \left( \frac{m_i T}{2\pi} \right)^{3/2} e^{-m_i/T}$$

☹️ Statistical independence at the classical level

☹️ Non-interacting

$$i = \{A_i, Z_i\}$$

😊 Analytical calculations

## □ This work

😊 Exact quantum counting

😊 Coulomb interaction + excluded volume

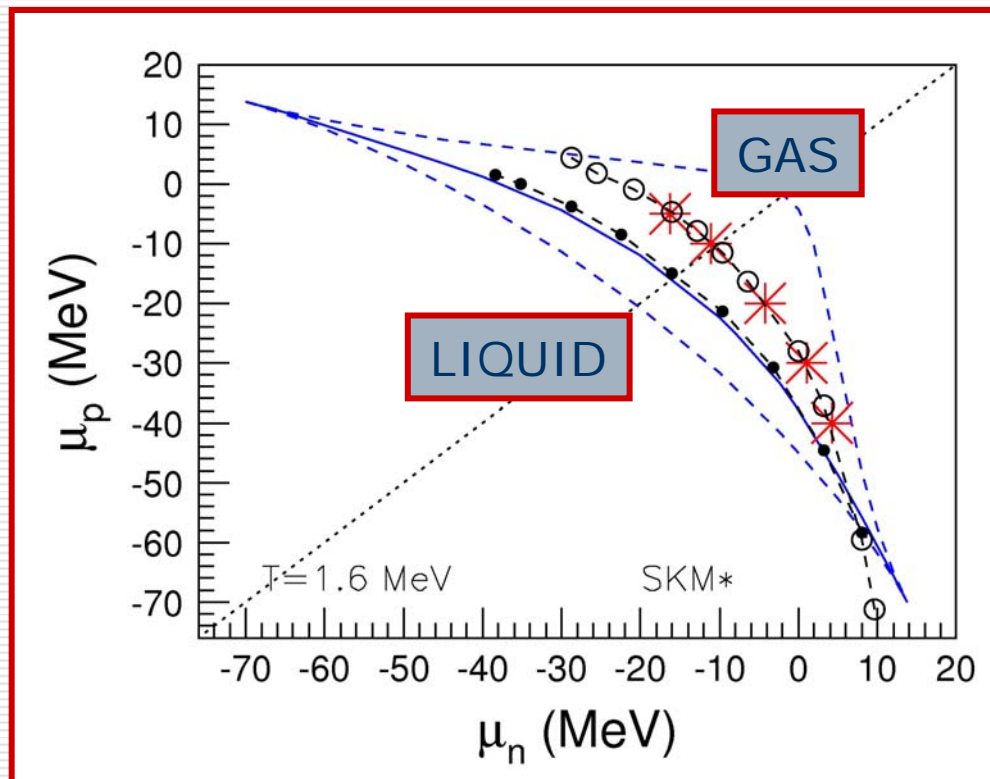
☹️ Expensive MC calculations

☹️ Convergence to be checked

$$Z = \sum_K \frac{(V - V_K)^{N_K} e^{-E_{coul}(K)/T}}{N_K!} \prod_{i=1}^{N_K} Z_i$$

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# Thermodynamics of clusterized and unclusterized matter

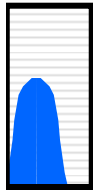


- Coexistence of purely homogeneous NM
- - - Spinodal of purely homogeneous NM
- \* Coexistence of purely clusterized NM
- ○ Trajectories of purely clusterized NM
- ● Trajectories of purely clusterized NM

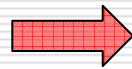
=> The mean-field instabilities are cured by cluster formation

# An hybrid model for the crust-core transition

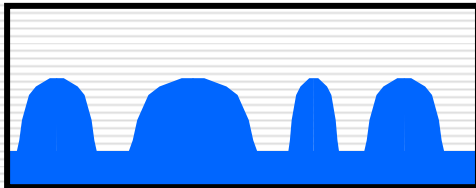
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Nuclei in the outer crust



Statistical ensemble of interacting excited clusters



Neutron (proton) drip in the inner crust → the two components together



Homogeneous matter in the core

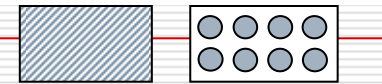


Finite temperature Hartree Fock with Skyrme interactions (Sly230a)

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+ electrons  $\rho_e = \rho_p$ ,  $\gamma, \nu$

# Phase mixture versus phase coexistence



A system composed of heterogenous components I=HM, II=clus

## □ Mixture

(ex: atmosphere)  
dishomogeneities on a  
microscopic scale

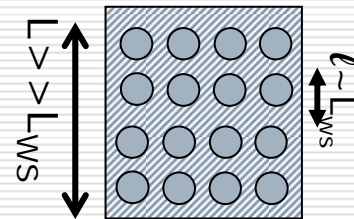
$$\beta^{HM} = \beta^{clus}$$

$$\mu_i^{HM} = \mu_i^{clus} \quad i = n, p$$

$$P = P^{HM} + P^{clus}$$

$$\rho_i = \rho_i^{HM} + \rho_i^{clus} \quad (+ \text{ excluded volume})$$

=> **Continuous EOS**



## □ Coexistence

(ex: Solid-Liquid)  
dishomogeneities on a  
macroscopic scale

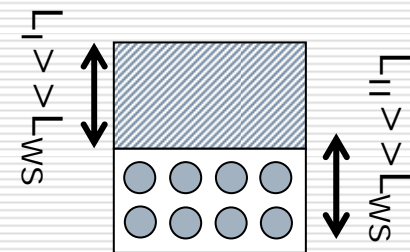
$$\beta^{HM} = \beta^{clus}$$

$$\mu_i^{HM} = \mu_i^{clus} \quad i = n, p$$

$$P^{HM} = P^{clus}$$

$$\rho_i = \rho_i^{HM} x + \rho_i^{clus} (1-x) \quad (\text{Gibbs construction})$$

=> **jump in observables**

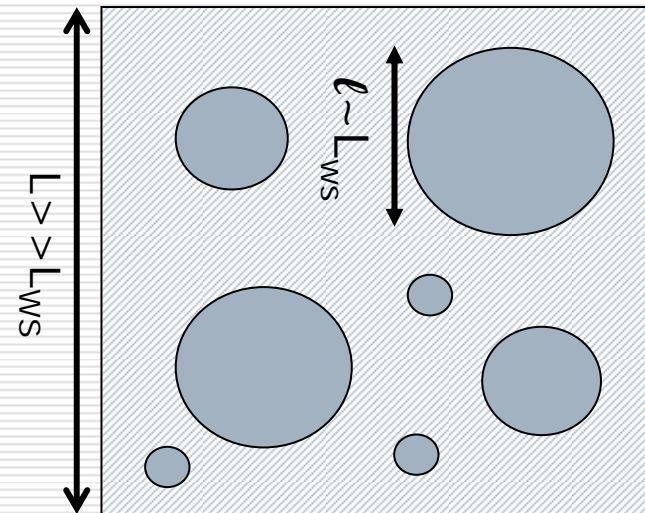


# No first order transition in dilute stellar matter

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A first order crust-core transition (e.g. Lattimer-Swesty, Shen, etc.)

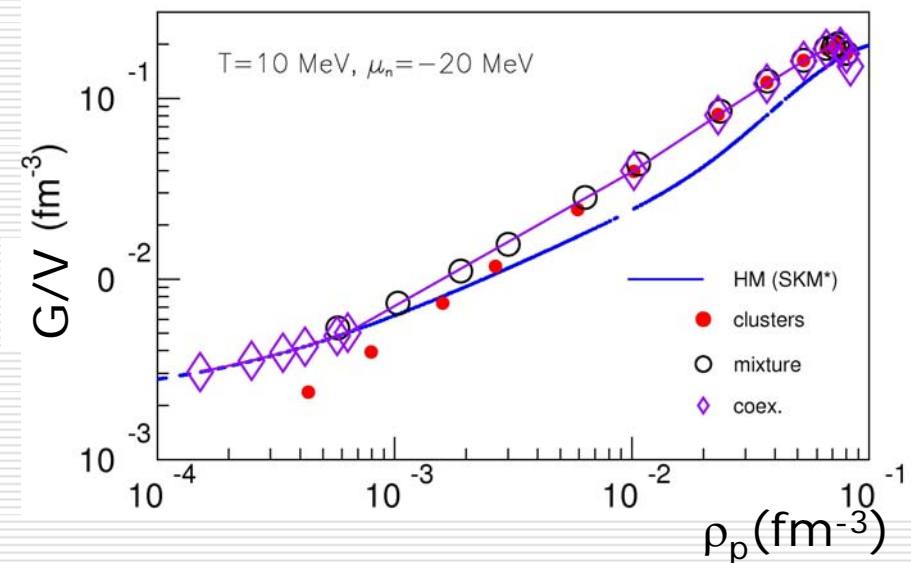
- Does not correspond to the physical structure of the crust (*microscopic fluctuations*)



# No first order transition in dilute stellar matter

A first order crust-core transition (e.g. Lattimer-Swesty, Shen, etc.)

- Does not correspond to the physical structure of the crust (*microscopic fluctuations*)
- Gives no entropy gain



# No first order transition in dilute stellar matter

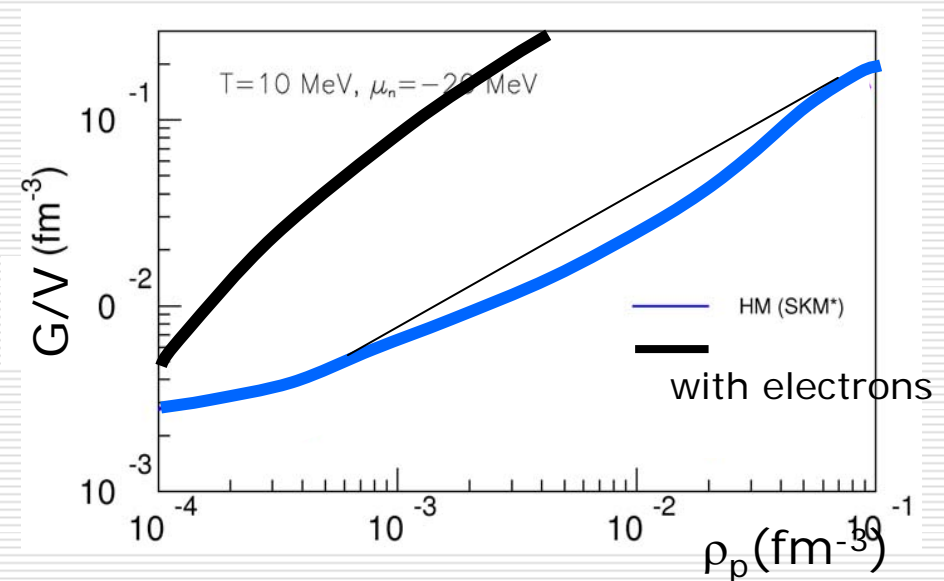
A first order crust-core transition (e.g. Lattimer-Swesty, Shen, etc.)

□ Does not correspond to the physical structure of the crust  
(*microscopic fluctuations*)

□ Gives no entropy gain

□ Ignores electron incompressibility!!!

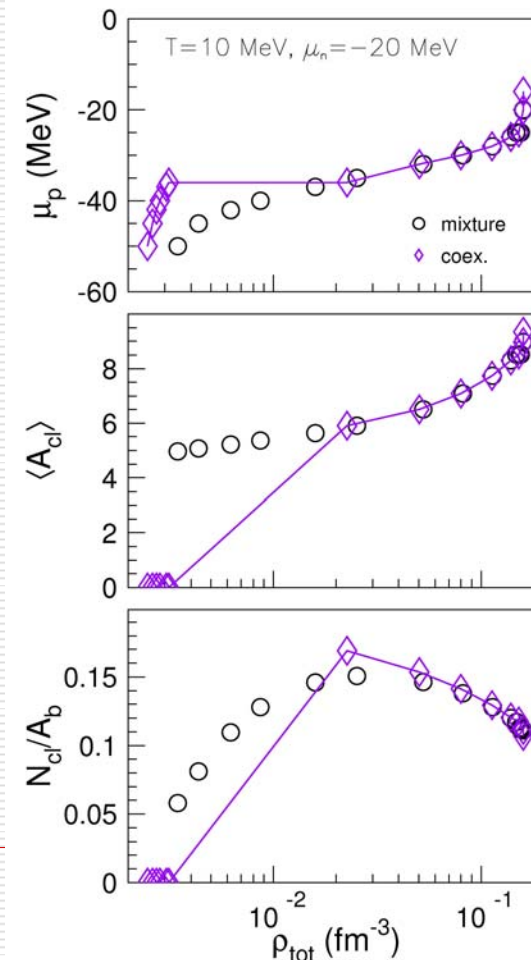
(*transition quenched because  $\partial\mu_e/\partial\rho_e \approx \text{GeV}$   
 $\Rightarrow$  concave entropy*)



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A first order crust-core transition (e.g. Lattimer-Swesty, Shen, etc.)

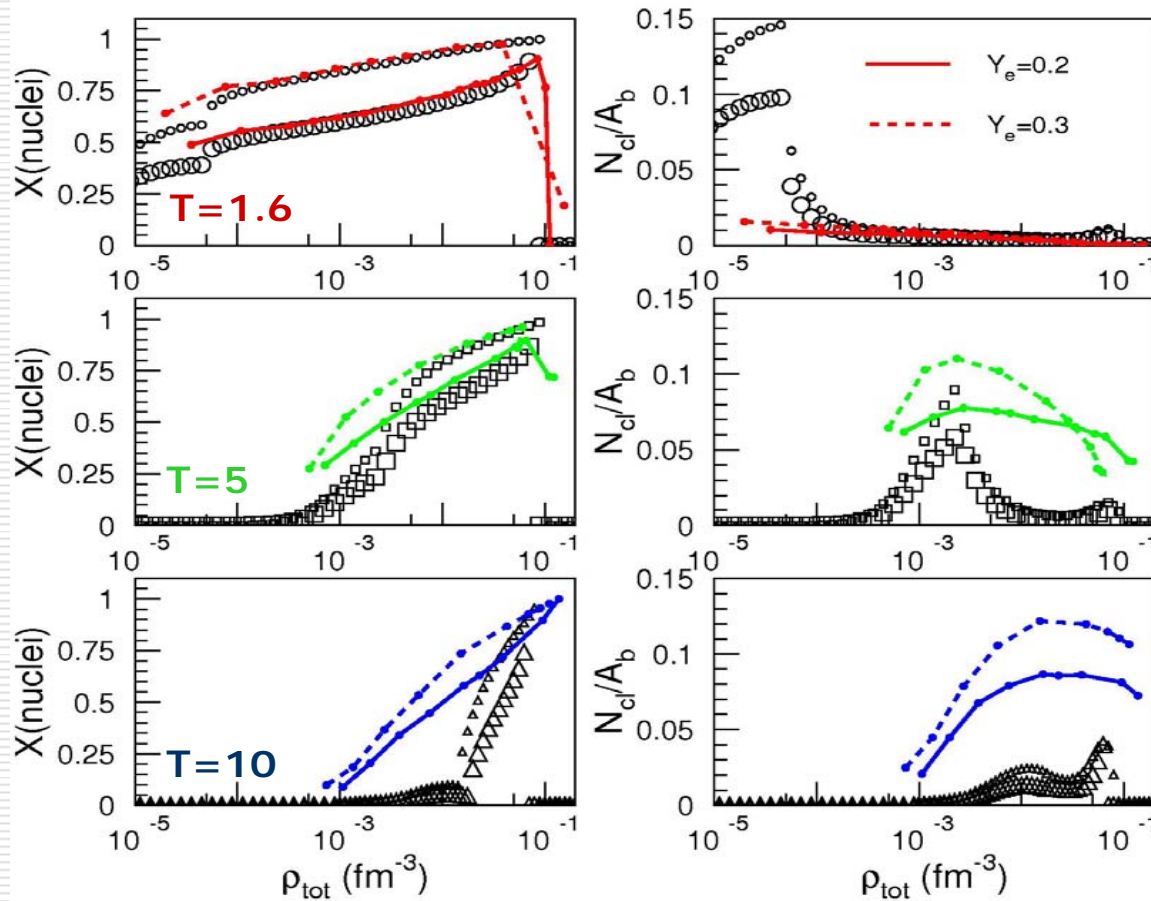
- ❑ Does not correspond to the physical structure of the crust (*microscopic fluctuations*)
- ❑ Gives no entropy gain
- ❑ Ignores electron incompressibility!!!  
(*transition quenched because  $\partial\mu_e/\partial\rho_e \approx \text{GeV}$   
 $\Rightarrow$  concave entropy*)
- ❑ Produces artificial discontinuities  
(*problem with SN codes*)





# Crust composition: cluster contribution

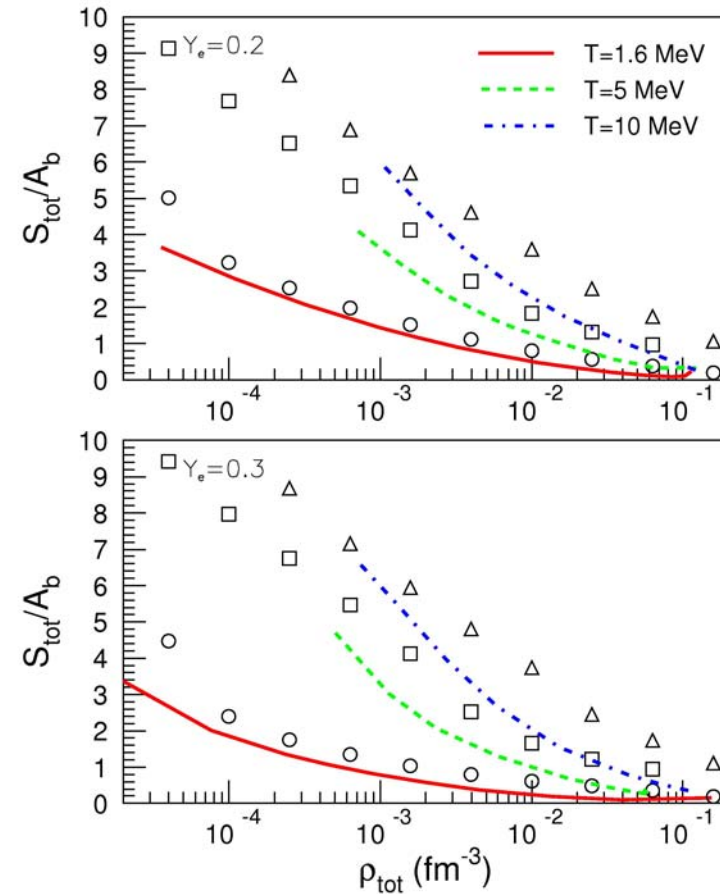
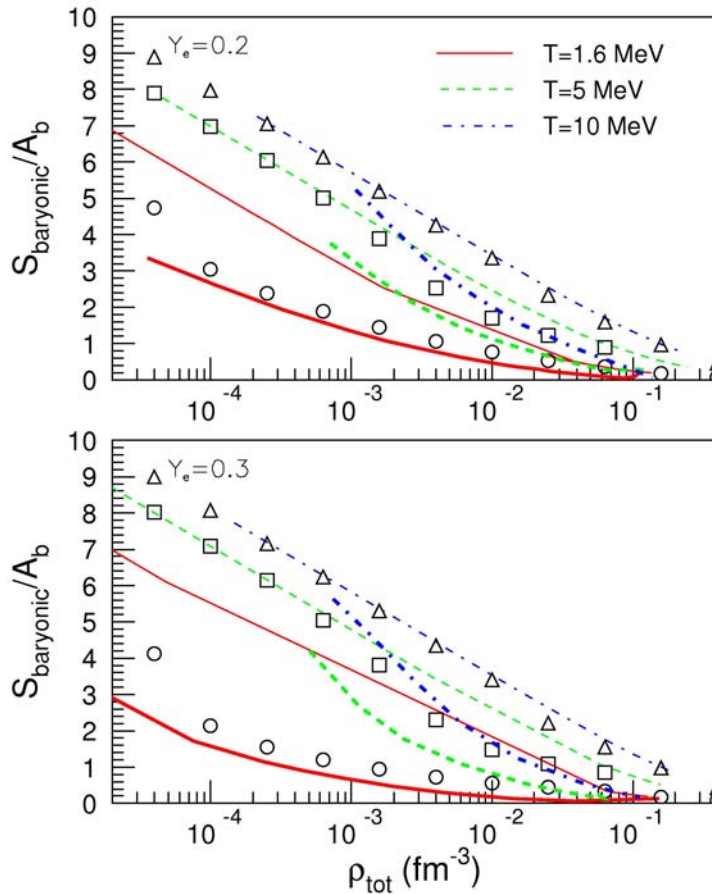
Lines: this work  
Symbols: LS EOS



- Decreasing cluster size with increasing temperature
- Clusters still important at  $T = 10$  MeV

# Entropy density

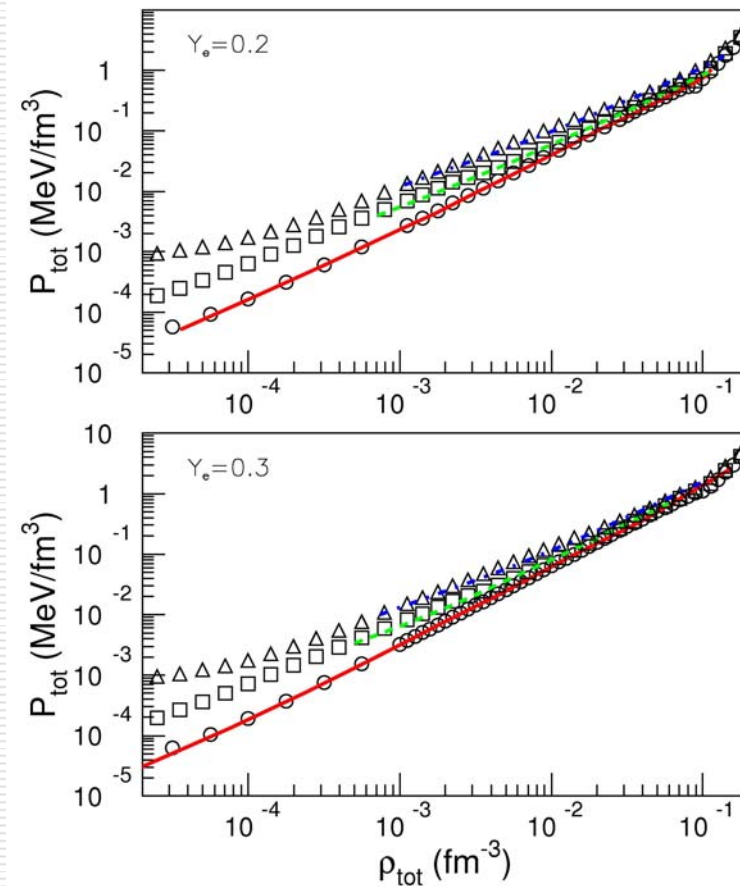
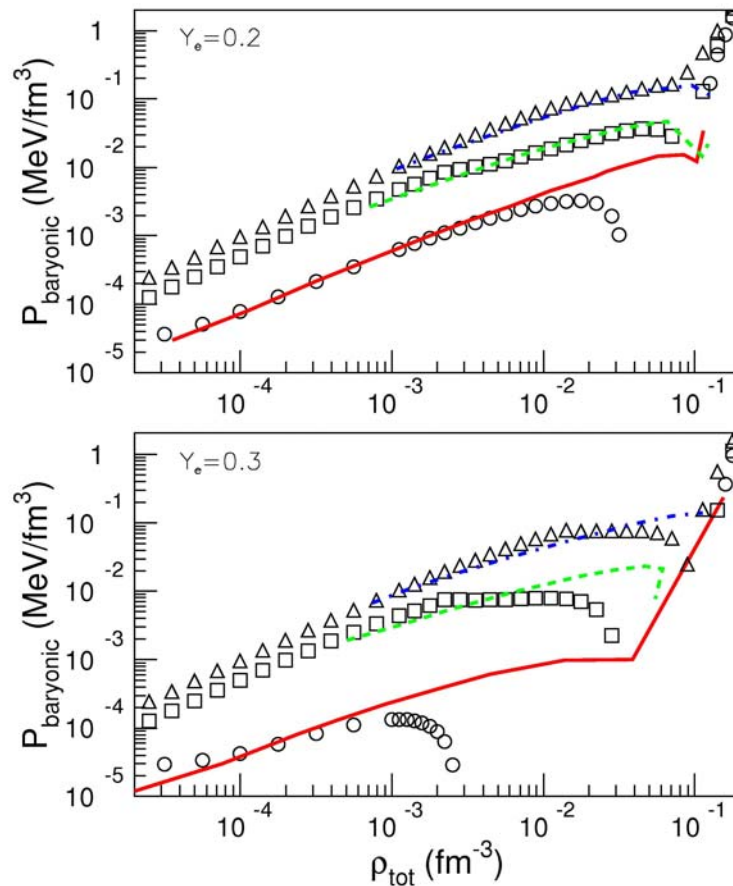
Symbols: LS EOS  
Thick Lines: this work  
Thin lines: clusters excluded



Differences with LS at high temperature even in the total entropy, due to the presence of clusters

# Pressure

Lines: this work  
Symbols: LS EOS



Differences with LS at high density, due to the absence of a first order transition

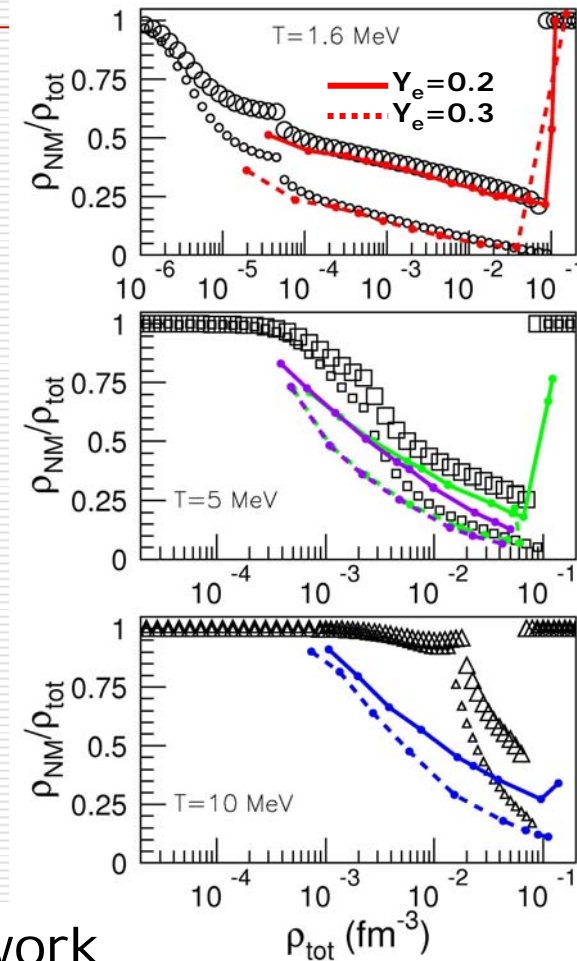
# Density and pressure at the crust-core transition

*J.M. Lattimer and M. Prakash, Phys. Rep. 442, 109 (2007).*

- Crustal fraction of the moment of inertia

$$\frac{\Delta I}{I} = \frac{28\pi P_t R^3}{3Mc^2 \xi} \frac{1 - 1.67\xi - 0.6\xi^2}{1 + \frac{2P_t(1 + 5\xi - 14\xi^2)}{\rho_t mc^2 \xi^2}}$$

- Can be measured from pulsar glitches
- Puts constraints on the NS radius; ex: Vela pulsar
- Transition naturally obtained !



Lines: this work  
Symbols: LS EOS

# Conclusion

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- ❑ Specific thermodynamics for the dilute matter in NS crusts and SN cores
  - ❑ Model-independent conclusion: no first-order phase transition
  - ❑ Illustration within an improved NSE model combining nuclear matter properties and all-sized clusters
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[Redacted]

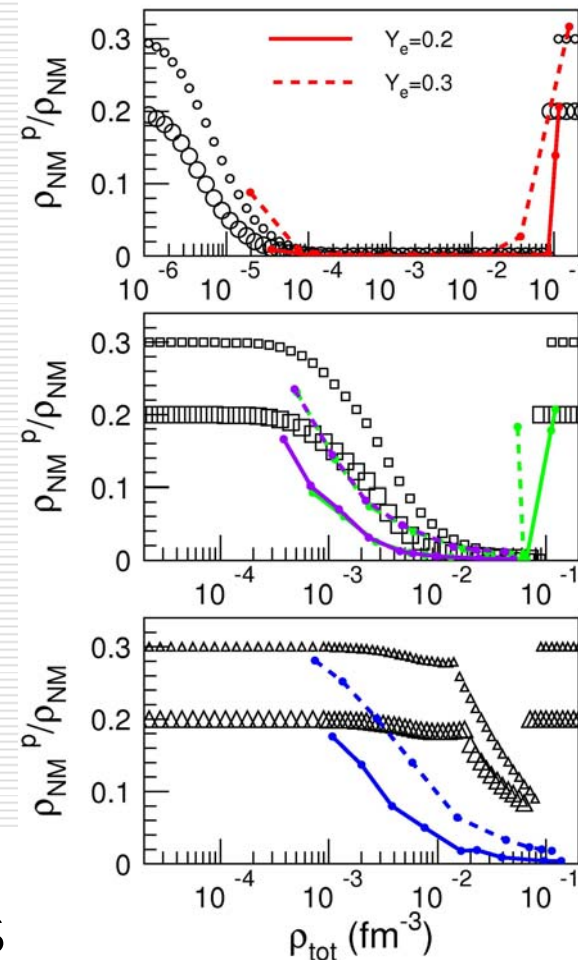
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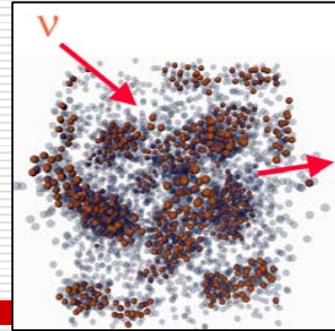
# Neutron versus proton drip

- Proton drip is negligible at low temperature
- Increases at high temperature, but much less than in LS

Lines: this work  
Symbols: LS EOS



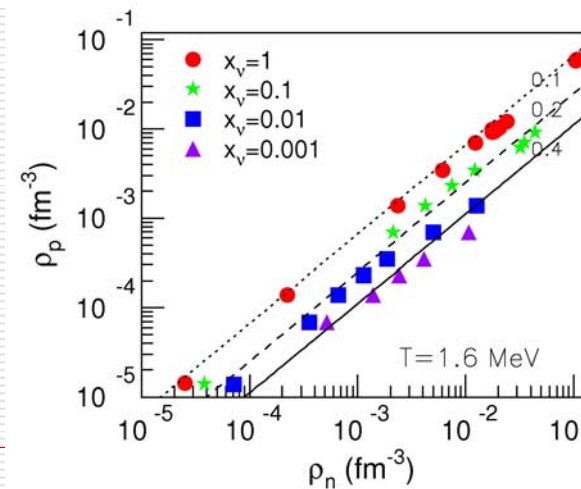
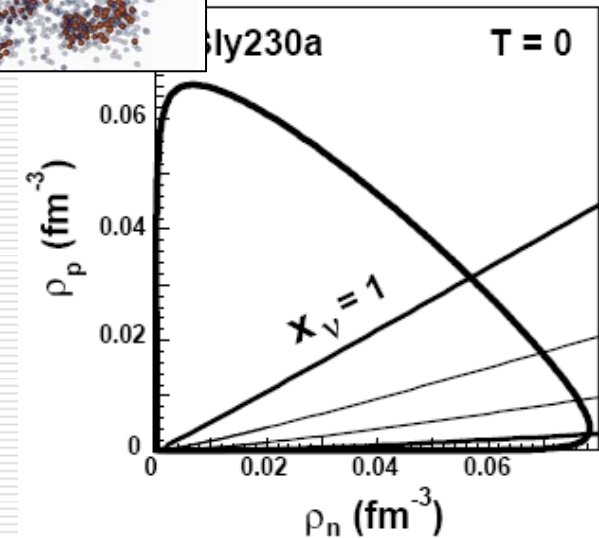
# Neutrino opacity



- $x_\nu$  percentage of trapped neutrinos in  $\beta$  equilibrium
- Determines the leptonization rate  $\Rightarrow$  the size of the homologous core
- Similar results to MF calculation with the same effective interaction

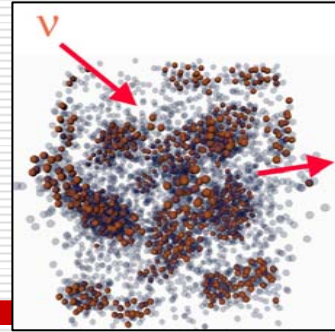
$$x_\nu = \frac{1}{2\pi^2 \hbar^3 \rho_\nu^{prod}} \int_0^{2\mu_\nu} de e^2 n_{\beta\mu}(e)$$

*C. Ducoin et al, NPA2007*

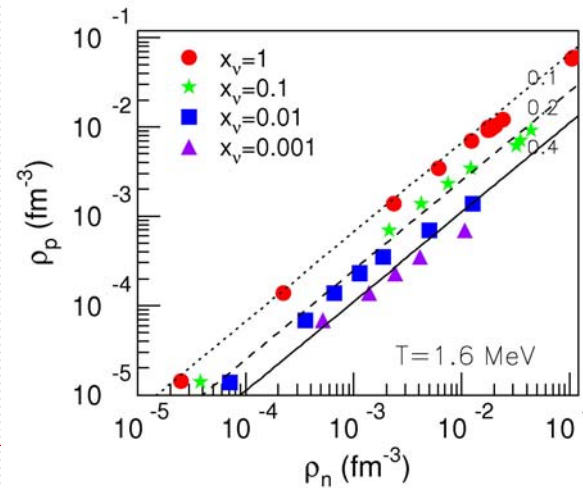
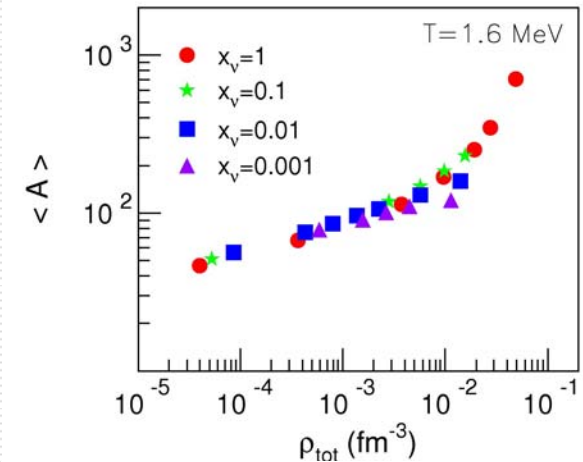




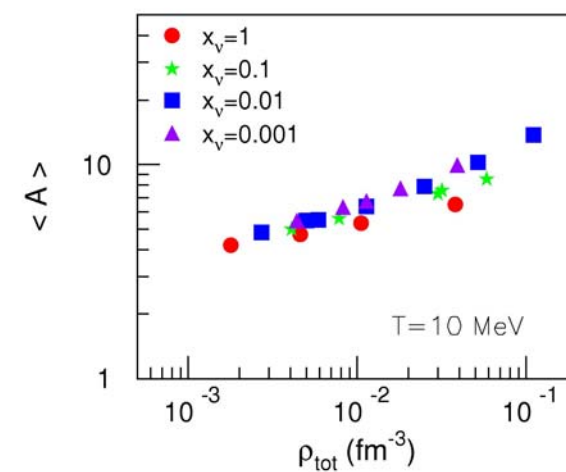
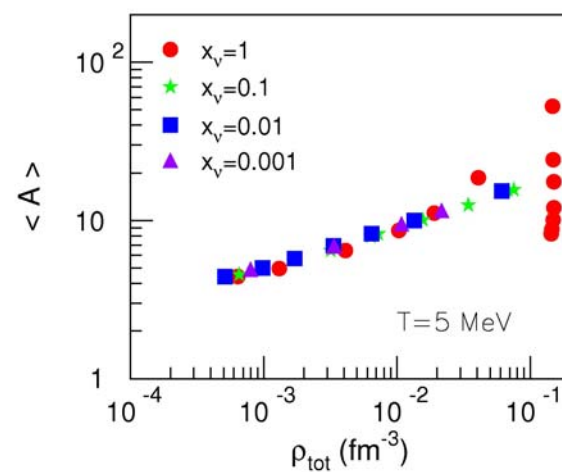
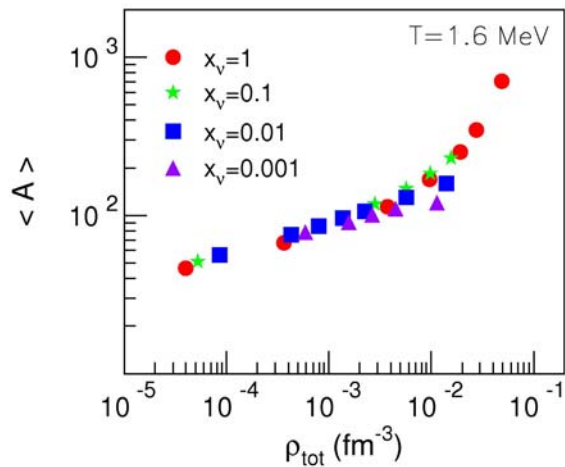
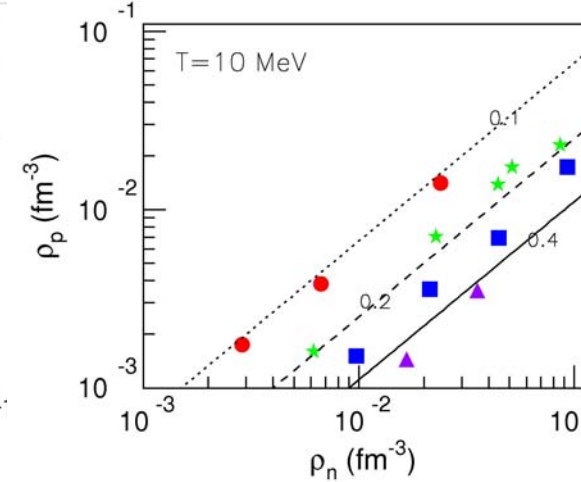
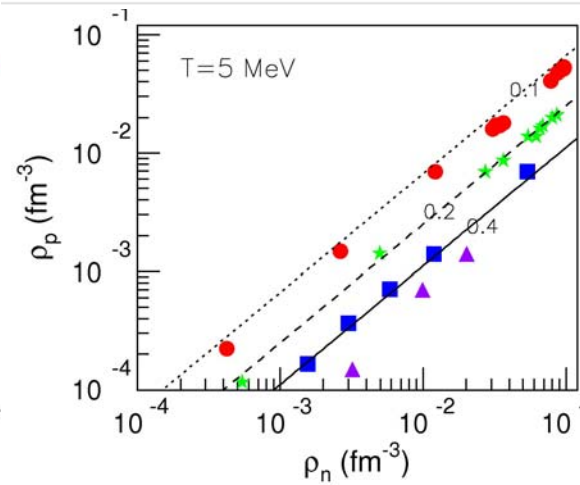
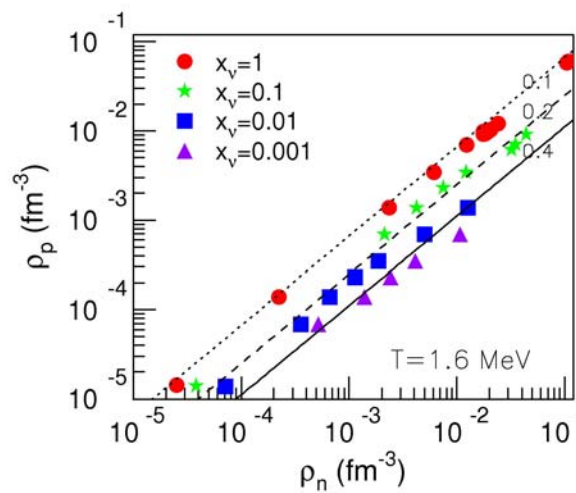
# Neutrino opacity



- $x_\nu$  percentage of trapped neutrinos
- Determines the leptonization rate  $\Rightarrow$  the size of the homologous core
- Similar results to MF calculation with the same effective interaction
- Can be correlated to the size of the clusters and number of free protons

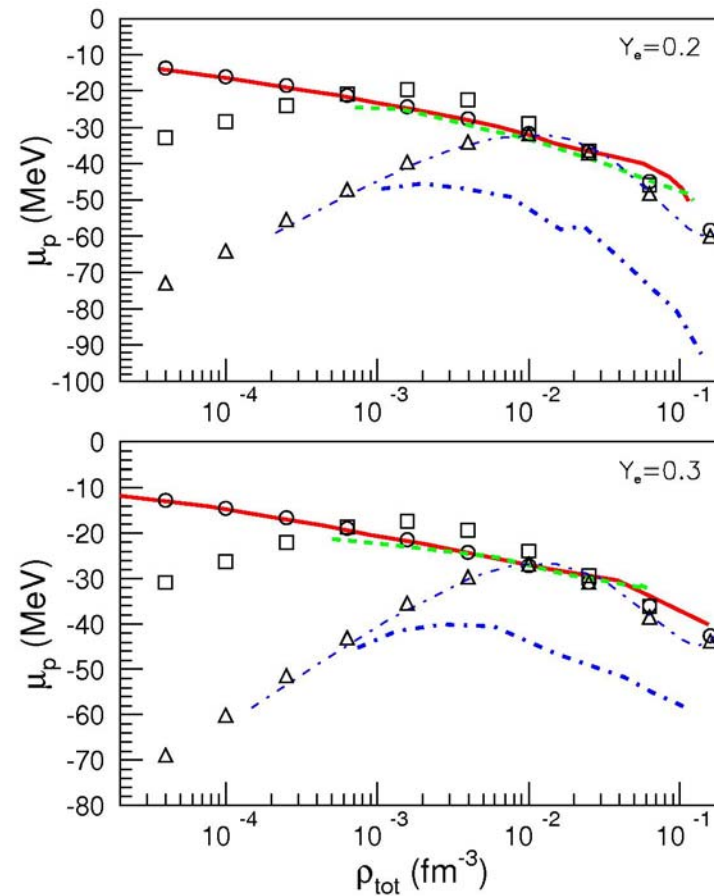
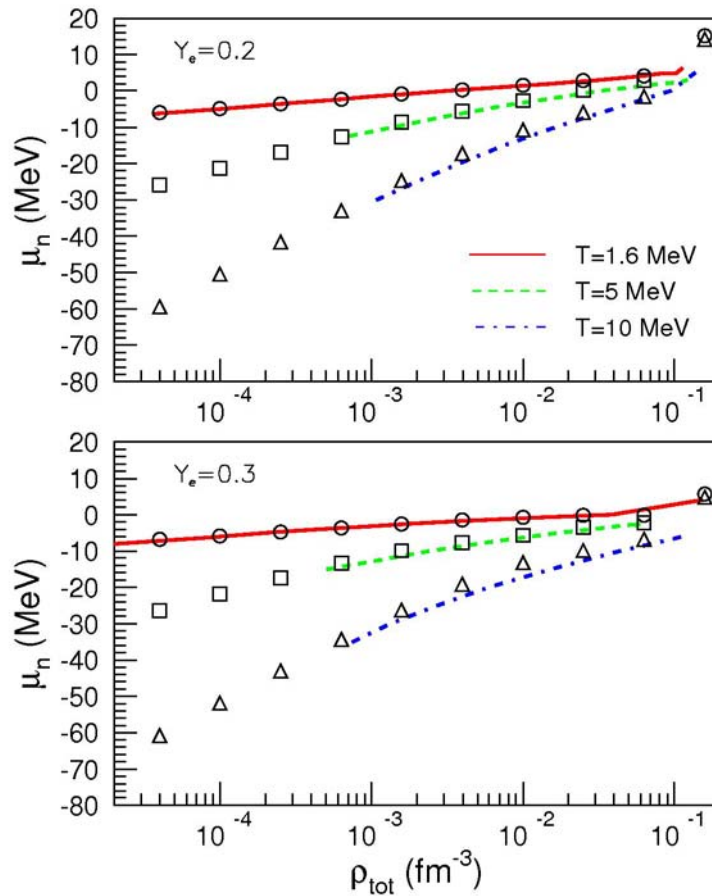


# Opacity to neutrinos



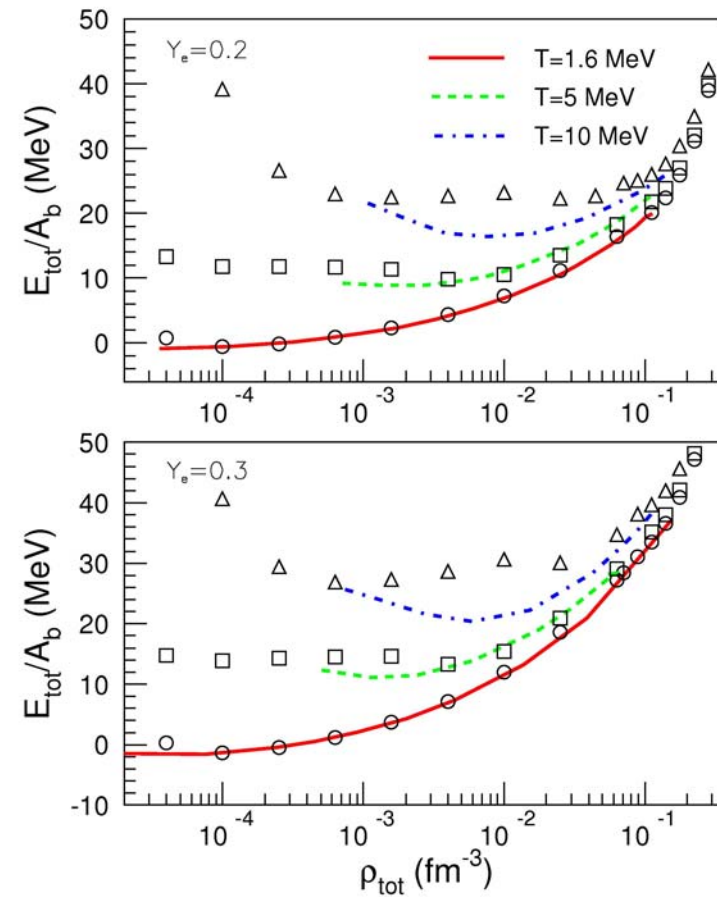
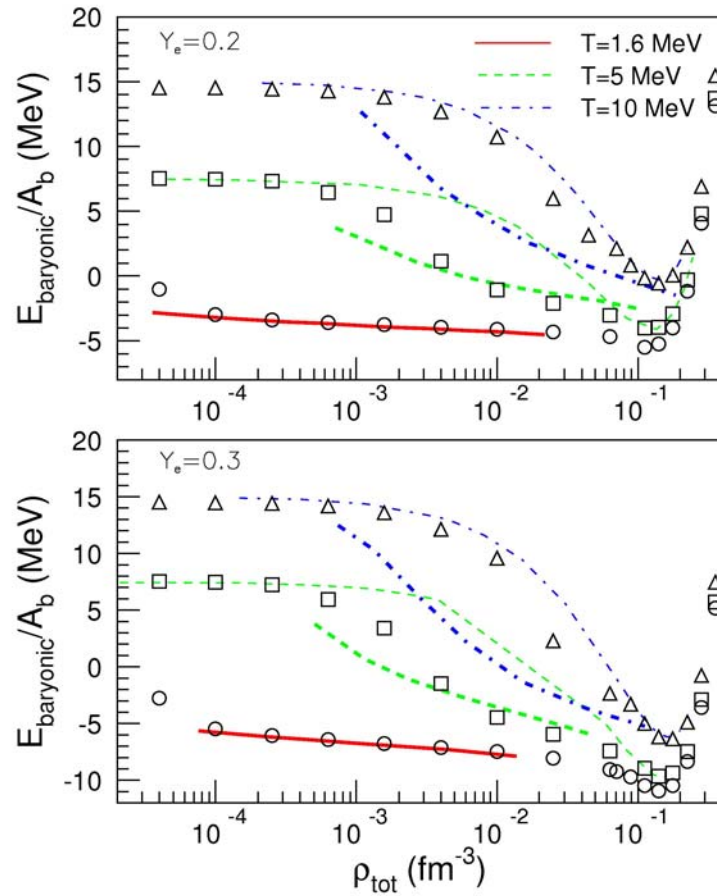
# Chemical potentials

Symbols: LS EOS  
Thick Lines: this work  
Thin lines: clusters excluded

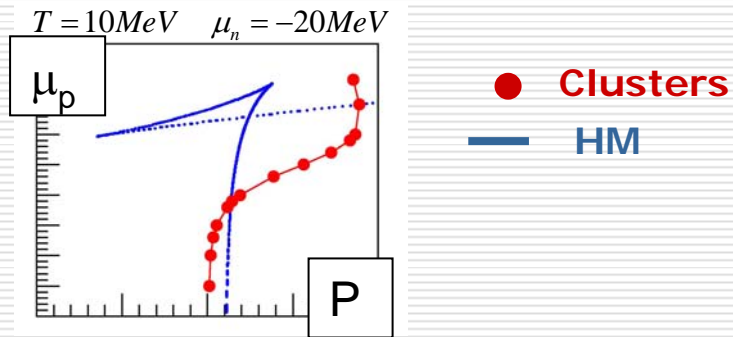


# Energy density

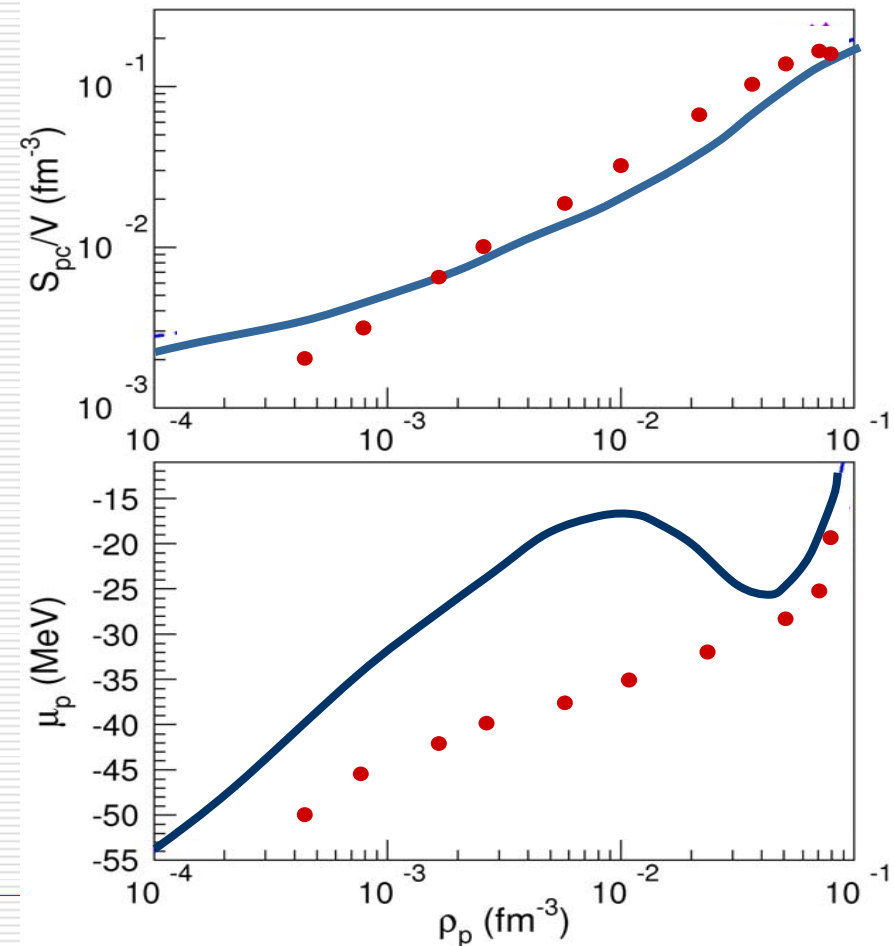
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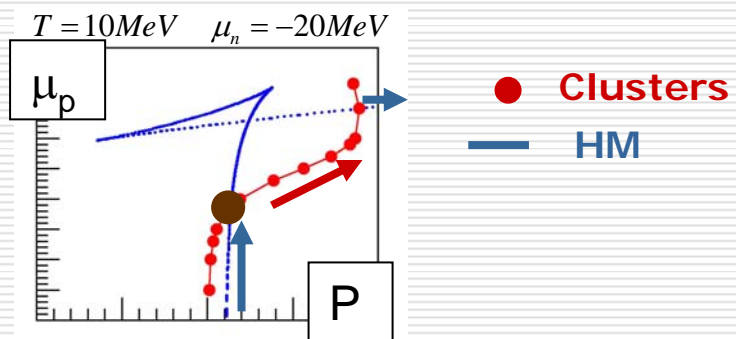
# The order of the crust-core transition



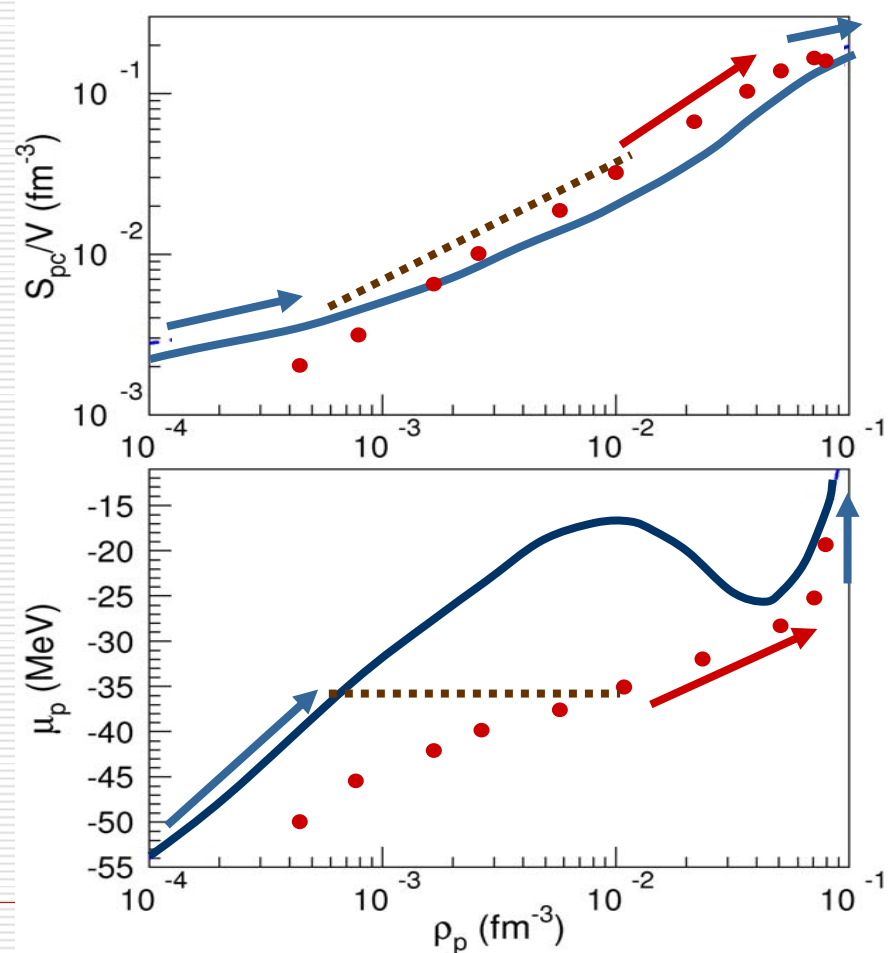
- First order  $\Phi$  transition:



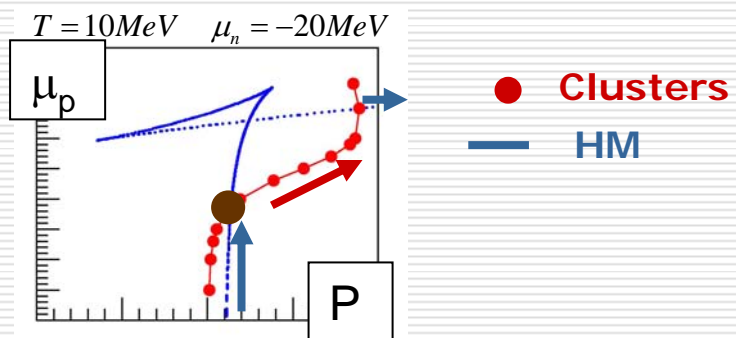
# The order of the crust-core transition



- First order  $\Phi$  transition:  
discontinuous change from **HM**  
Through a coexistence point  
To **clusterized matter**  
To **Homogeneous Matter** again



# The order of the crust-core transition



- First order  $\Phi$  transition: discontinuous change from **HM** Through a coexistence point To **clusterized matter** To **Homogeneous Matter** again

- Mixture: ● Continuous EOS maximising the same entropy

