

# **Neutron-proton asymmetry in nuclear matter and finite nuclei**

***Dao Tien Khoa***

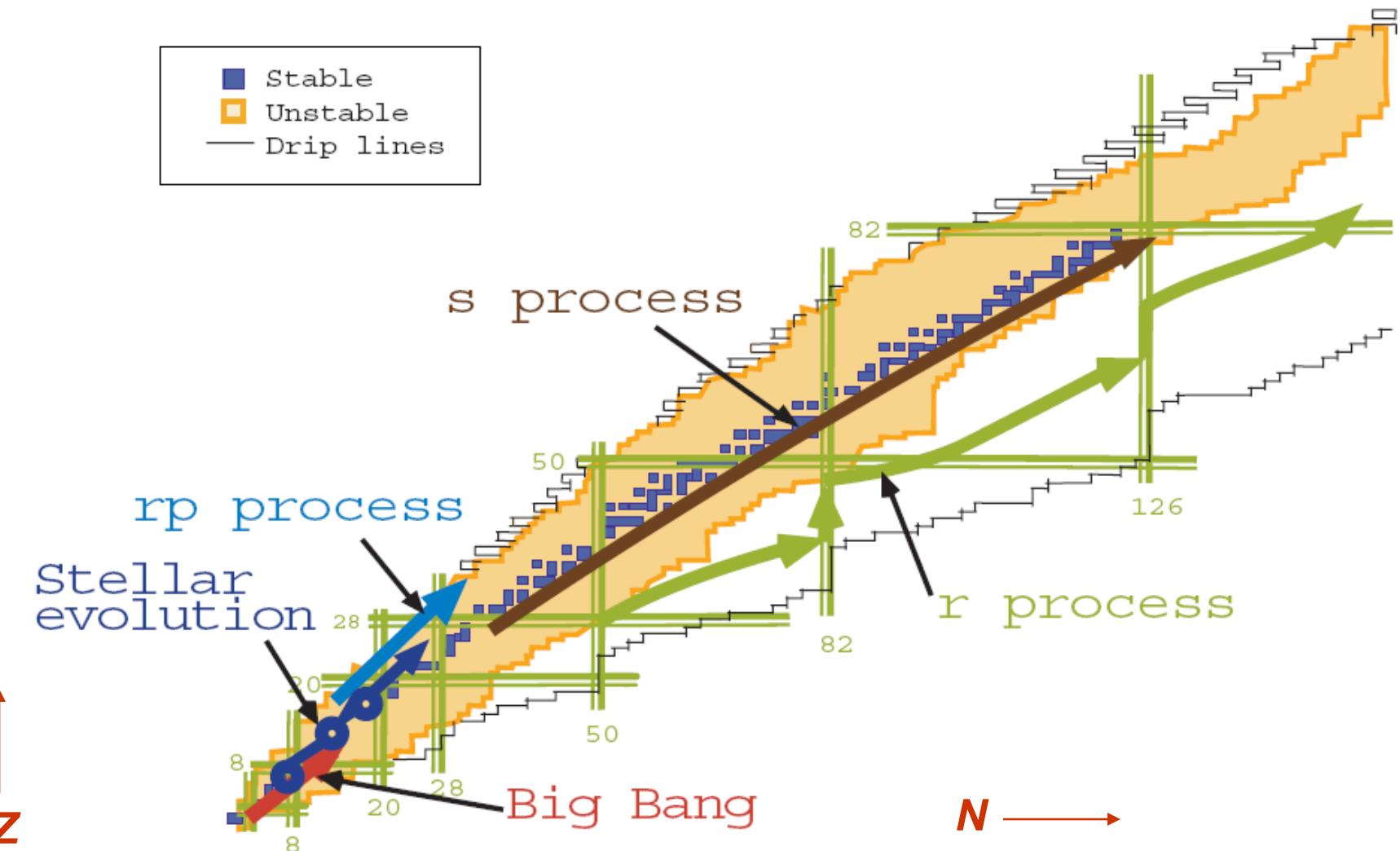
*Institute for Nuclear Science & Technology  
Vietnam Atomic Energy Commission (VAEC)*

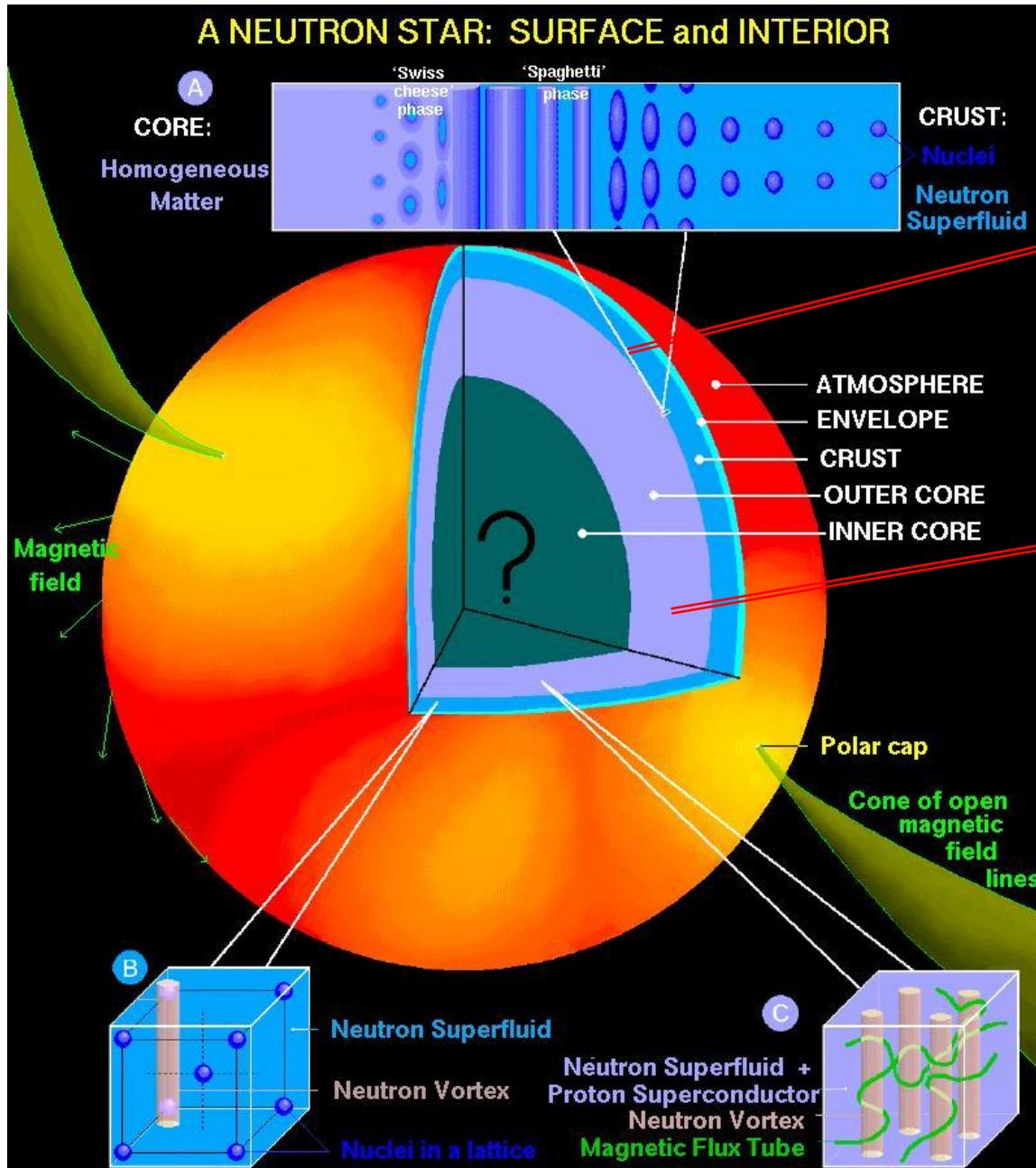
- *Equation of state for asymmetric nuclear matter*
- *Charge-exchange ( $p,n$ )IAS reaction*  $\iff$  *Nuclear symmetry energy*
- *( $p,p'$ ) scattering on Oxygens*  $\iff$  *Impurity of isospin symmetry ?*

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# Neutron-proton asymmetry in finite nuclei $\delta = (N-Z)/A$

$\delta$  is large in neutron-rich nuclei, with  $\delta_{\text{max}}=0.5$  for  ${}^8\text{He}$  !





Proton fraction  
 $x = \rho_p / \rho = 0.5 * (1 - \delta)$

$$\rho = (0.5 \sim 1) \rho_0$$

$$\delta = 0.94 \sim 0.90$$

$$x = 0.03 \sim 0.05$$

$$\rho_0 \sim 0.17 \text{ nucleon/fm}^{-3}$$

$$\rho = (2 \sim 6) \rho_0$$

$$\delta = 0.86 \sim 0.80$$

$$x = 0.07 \sim 0.10$$

Sly EOS by Douchin & Haensel  
*Astronomy & Astrophysics*  
**380** (2001) 151

Experimentally

$\delta \leftrightarrow$  Symmetry Energy  
still unknown at large  $\rho$  !

# **Microscopic calculation of nuclear matter**

BHF or DBHF

Talks by Schulze & Lombardo



$$G[\omega; \rho] = V + \sum_{k_a, k_b > k_F} V \frac{|k_a k_b\rangle \langle k_a k_b|}{\omega - e(k_a) - e(k_b) + i\epsilon} G[\omega; \rho]$$

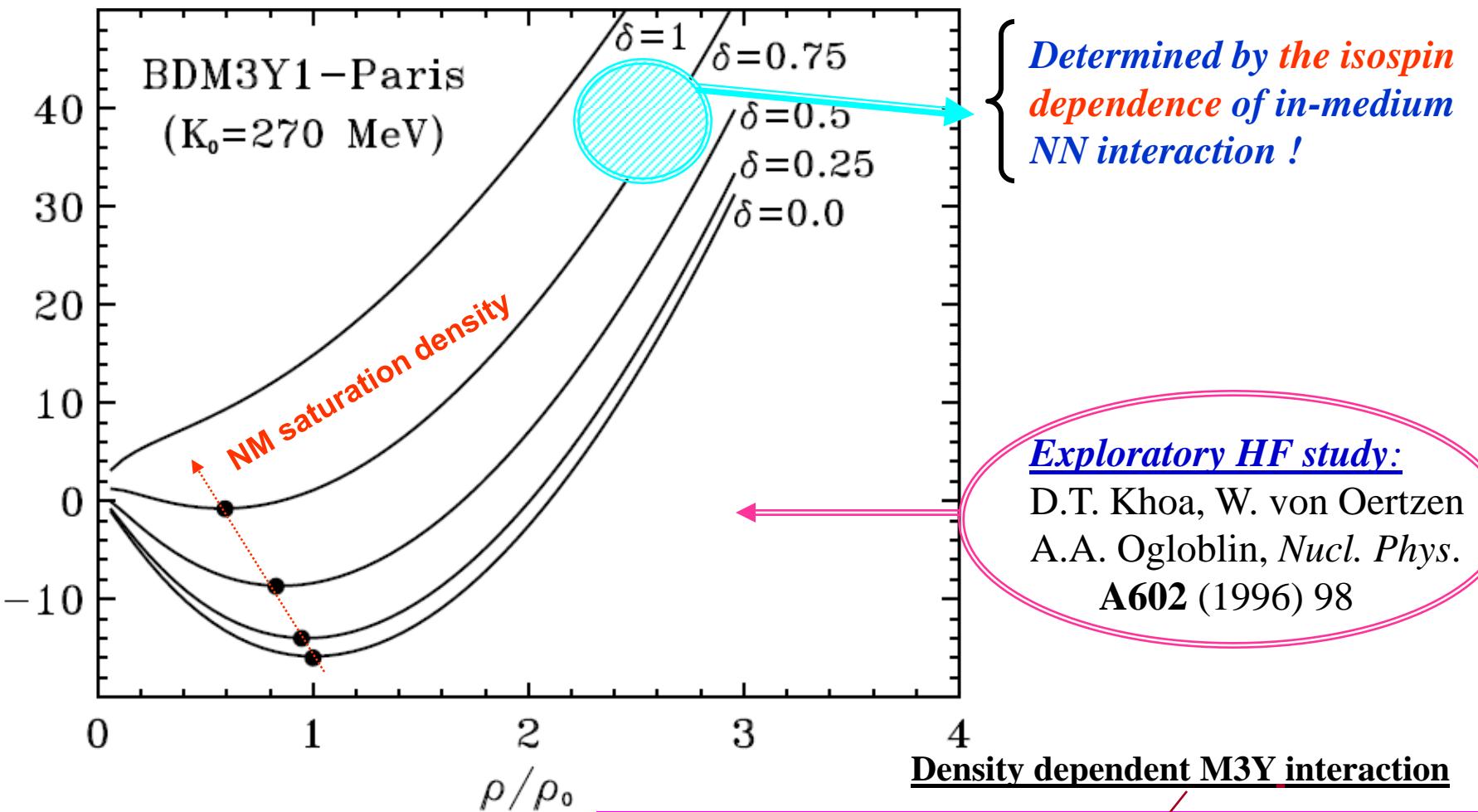
$$\frac{B}{A} = \frac{3}{5} \frac{k_F^2}{2m} + \frac{1}{2\rho} \text{Re} \sum_{k, k' \leq k_F} \langle kk' | G[e(k) + e(k'); \rho] | kk' \rangle_a$$

Antisymmetrization

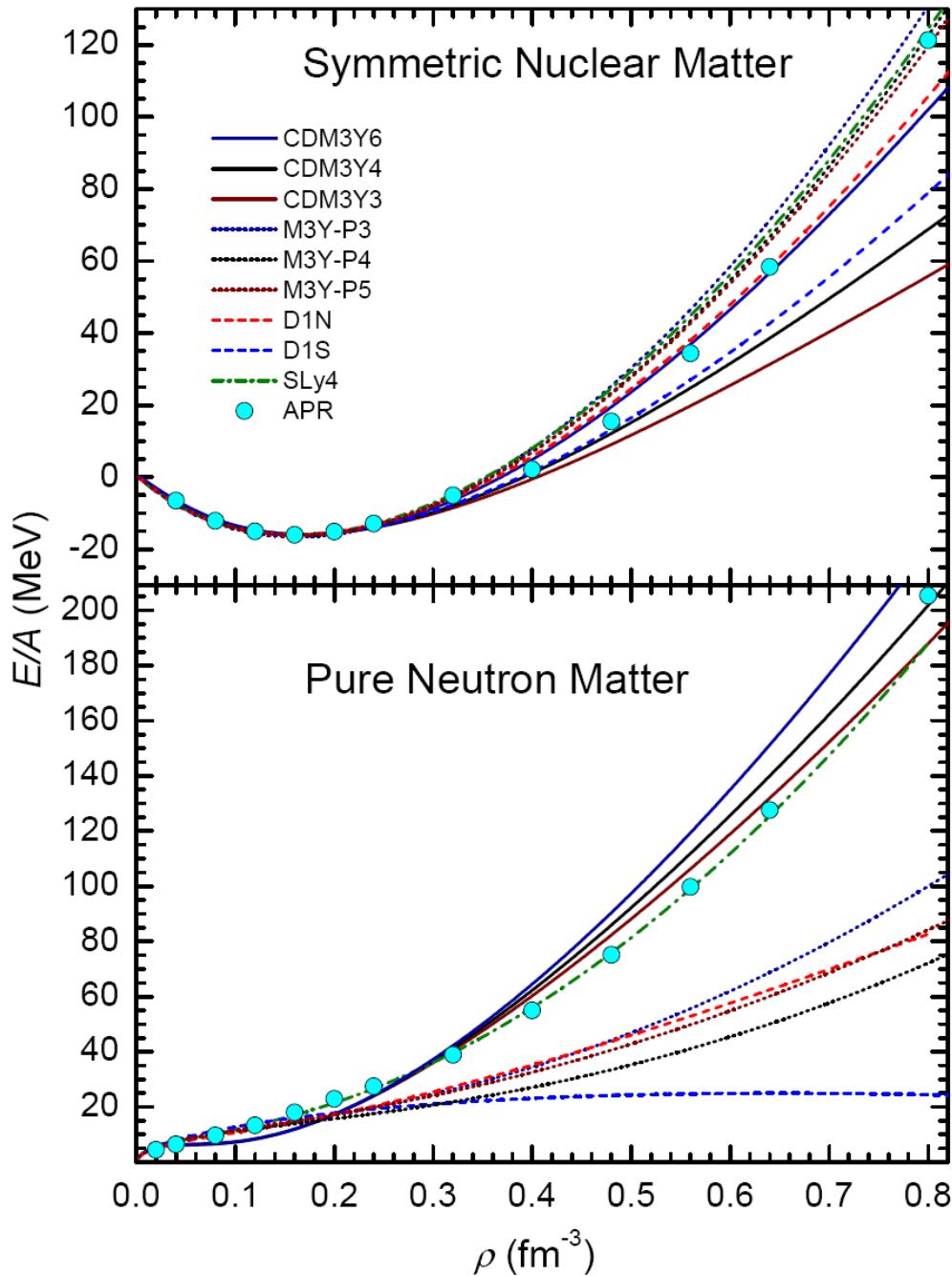
$$\frac{B}{A} = \frac{E}{A}(\rho, \delta) = \frac{E}{A}(\rho, \delta = 0) + S(\rho)\delta^2 + O(\delta^4) + \dots$$

Nuclear matter symmetry energy ( $E_{\text{sym}}$ )

## EOS of asymmetric nuclear matter



$$\frac{B}{A} = \frac{3}{5} \frac{k_F^2}{2m} + \frac{1}{2\rho} \operatorname{Re} \sum_{k,k' \leq k_F} \langle kk' | G[e(k) + e(k'); \rho] | kk' \rangle_a$$



$$\frac{E}{A}(\rho, \delta = 0) + S(\rho)\delta^2 + O(\delta^4) + \dots$$

*HF results given by some mean-field interaction*

**CDM3Yn:** D.T. Khoa, G.R. Satchler, and W. von Oertzen, *Phys. Rev. C* **56**, 954 (1997); D.T. Khoa, H.S. Than, and D.C. Cuong, *Phys. Rev. C* **76**, 014603 (2007).

**M3Y-Pn:** H. Nakada, *Phys. Rev. C* **78**, 054301 (2008).

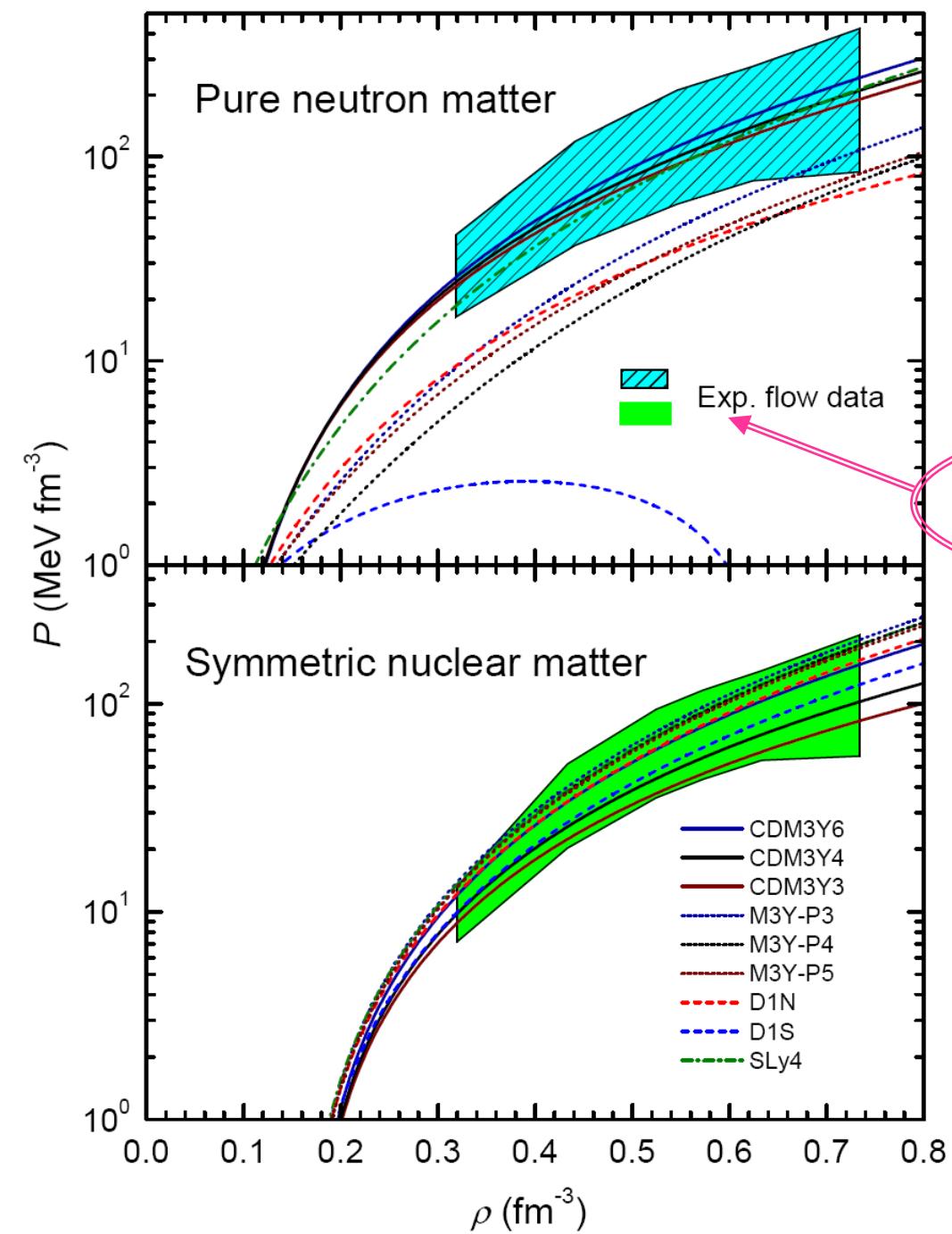
**D1S:** J.F. Berger, M. Girod, and D. Gogny, *Comp. Phys. Comm.* **63**, 365 (1991).

**D1N:** F. Chappert, M. Girod, and S. Hilaire, *Phys. Lett. B* **668**, 420 (2008).

**SLy4:** E. Chabanat *et al.*, *Nucl. Phys. A* **635**, 231 (1998)

*Ab-initio variational calculation using Argon V18 NN + NNN inter.*

**APR:** A. Akmal, V.R. Pandharipande, and D.G. Ravenhall, *Phys. Rev. C* **58**, 1804 (1998)



$$P(\rho, \delta) = \rho^2 \frac{\partial}{\partial \rho} \left[ \frac{E}{A}(\rho, \delta) \right]$$

P. Danielewicz, R. Lacey and W.G. Lynch  
 Science **298**, 1592 (2002)

M3Y-Pn, D1S, D1N *fail to reproduce empirical pressure of neutron matter !*

Two distinct scenarios for  
NM symmetry energy:  
*Asy-soft & Asy-stiff*



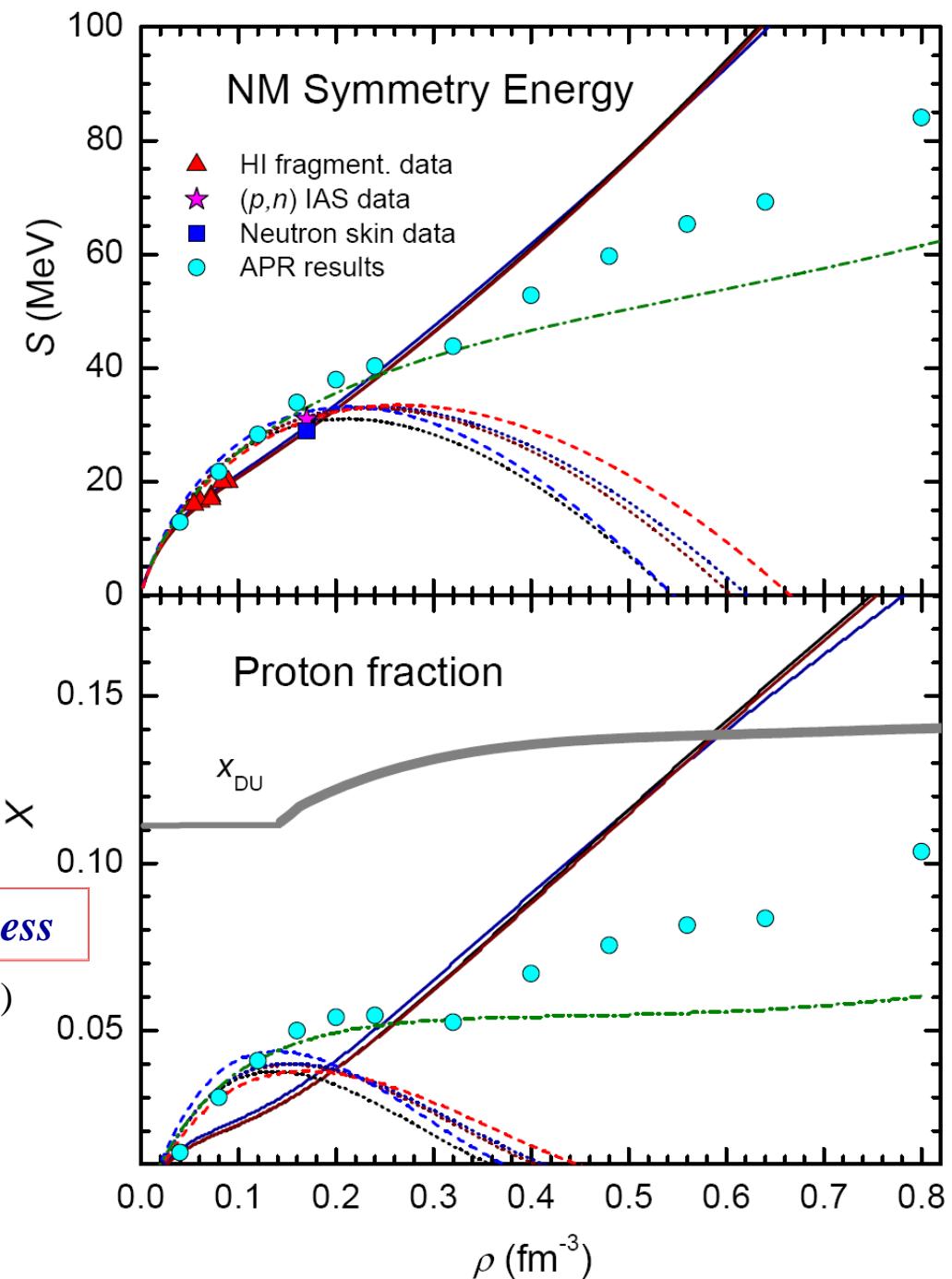
**Neutron star cooling ?**

H.S. Than, D.T. Khoa, N.V. Giai,  
*Phys. Rev. C* **80**, 064312 (2009).

**$x_{DU} \Rightarrow$  threshold for the direct Urca process**

T. Klahn *et al.*, *Phys. Rev. C* **74**, 035802 (2006)

**Talk by N.V. Giai !**



## *(p,n) charge exchange “scattering” to IAS states and isospin dependence of the nucleon optical potential*

Explicit **IS** and **IV** parts of the proton–nucleus potential

$$U(E, \mathbf{R}) = U_0(E, \mathbf{R}) - \varepsilon U_1(E, \mathbf{R}), \quad \varepsilon = \frac{N-Z}{A} = \frac{2T_A}{A}$$

$$U_0(E, \mathbf{R}) = \int \{ [\rho_p(\mathbf{r}) + \rho_n(\mathbf{r})] v_{00}^D(\rho, E, s) + \\ [\rho_p(\mathbf{R}, \mathbf{r}) + \rho_n(\mathbf{R}, \mathbf{r})] v_{00}^{EX}(\rho, E, s) j_0(k(E, R)s) \} d\mathbf{r},$$

$$U_1(E, \mathbf{R}) = \frac{1}{\varepsilon} \int \{ [\rho_n(\mathbf{r}) - \rho_p(\mathbf{r})] v_{01}^D(\rho, E, s) + \\ [\rho_n(\mathbf{R}, \mathbf{r}) - \rho_p(\mathbf{R}, \mathbf{r})] v_{01}^{EX}(\rho, E, s) j_0(k(E, R)s) \} d\mathbf{r}.$$

$U_1 \Rightarrow$  **microscopic description of Lane potential**

D.T. Khoa, E. Khan, G. Colo and N. van Giai, *Nucl. Phys.* **A706**, 61 (2002)

G. R. Satchler *et al.*, *Phys. Rev.* **136**, B637 (1964).

The explicit isospin coupling based on the total wave function

$$\Psi(\mathbf{R}) = |pA\rangle \chi_{pA}(\mathbf{R}) + |n\tilde{A}_{IAS}\rangle \chi_{n\tilde{A}_{IAS}}(\mathbf{R})$$

**=> the coupled channels equations for quasi-elastic ( $p, n$ ) scattering**

$$\left[ K_p + U_0(\mathbf{R}) - \frac{2T_A}{A} U_1(\mathbf{R}) + V_c(\mathbf{R}) - E_p \right] \chi_{pA}(\mathbf{R}) = -\frac{2\sqrt{2T_A}}{A} U_1(\mathbf{R}) \chi_{n\tilde{A}_{IAS}}(\mathbf{R})$$

$$\left[ K_n + U_0(\mathbf{R}) + \frac{2(T_A - 1)}{A} U_1(\mathbf{R}) - E_n \right] \chi_{n\tilde{A}_{IAS}}(\mathbf{R}) = -\frac{2\sqrt{2T_A}}{A} U_1(\mathbf{R}) \chi_{pA}(\mathbf{R})$$

$K_{p(n)}$  and  $E_{p(n)}$  are the kinetic-energy operators and center-of-mass energies of the entrance-channel and the exit-channel

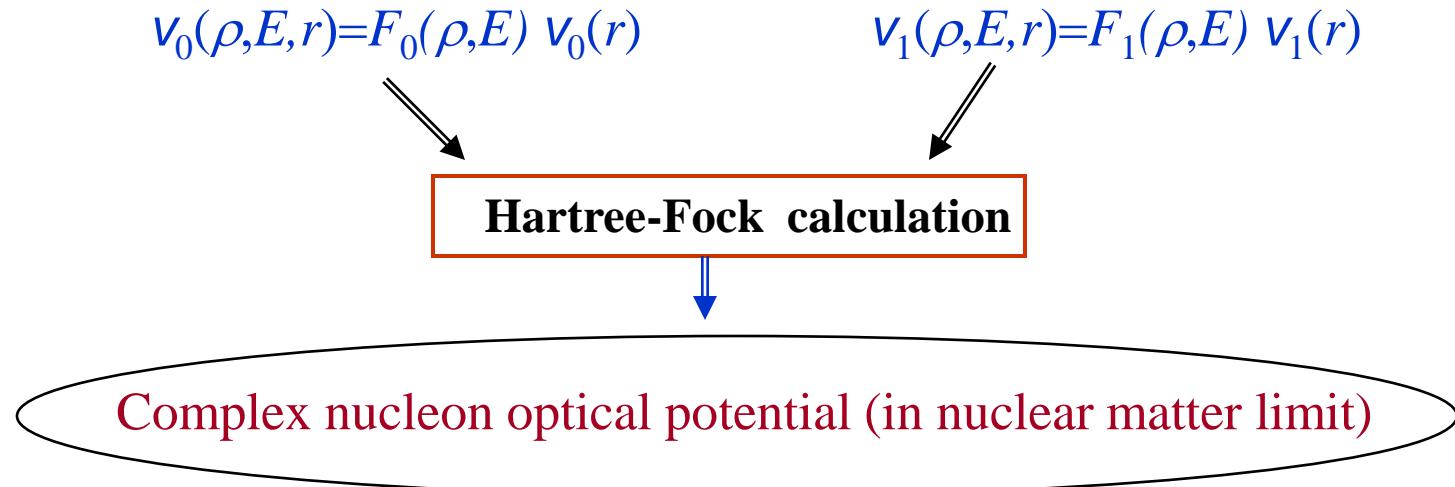
$$U_p(\mathbf{R}) = U_o(\mathbf{R}) - \frac{2T_A}{A} U_1(\mathbf{R}) \quad \Longrightarrow \quad \text{Central OP in the entrance channel}$$

$$U_n(\mathbf{R}) = U_0(\mathbf{R}) + \frac{2(T_A - 1)}{A} U_1(\mathbf{R}) \quad \Longrightarrow \quad \text{Central OP in the exit channel}$$

Folding model

Density- and isospin dependent NN interaction

**STEP I: *Mapping the isovector density dependence of the CDM3Yn interactions to the BHF results by JLM group***



$$U(\rho, E) = V_0(\rho, E) + i * W_0(\rho, E) \pm \delta * [V_1(\rho, E) + i * W_1(\rho, E)]$$

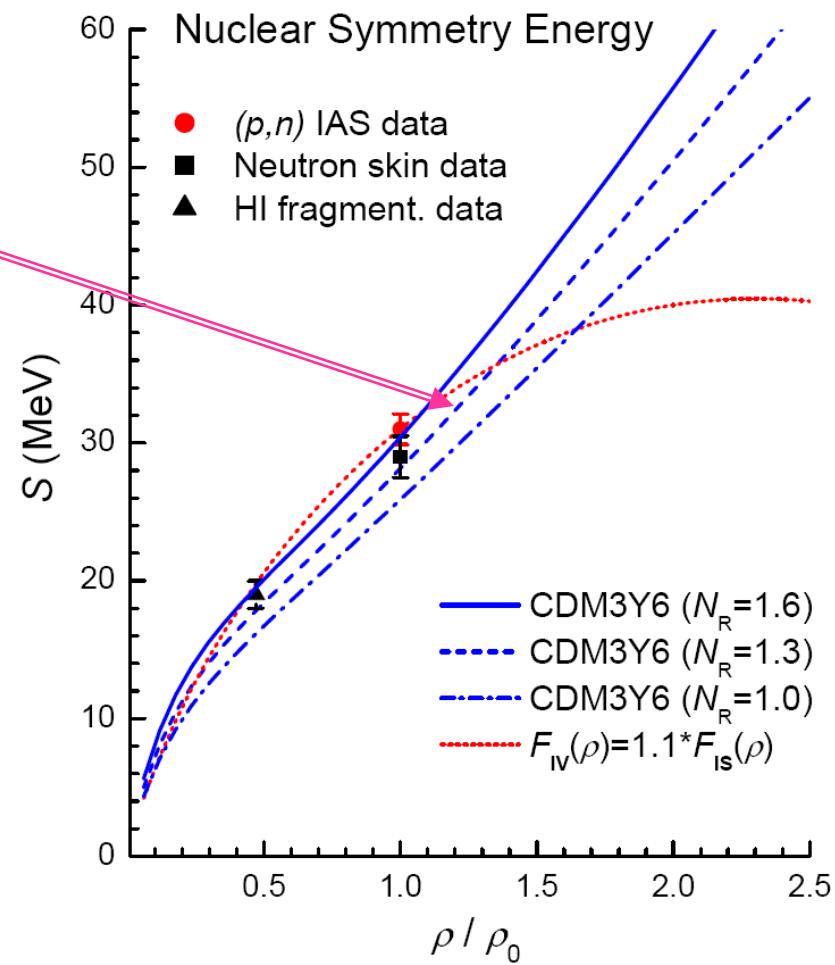
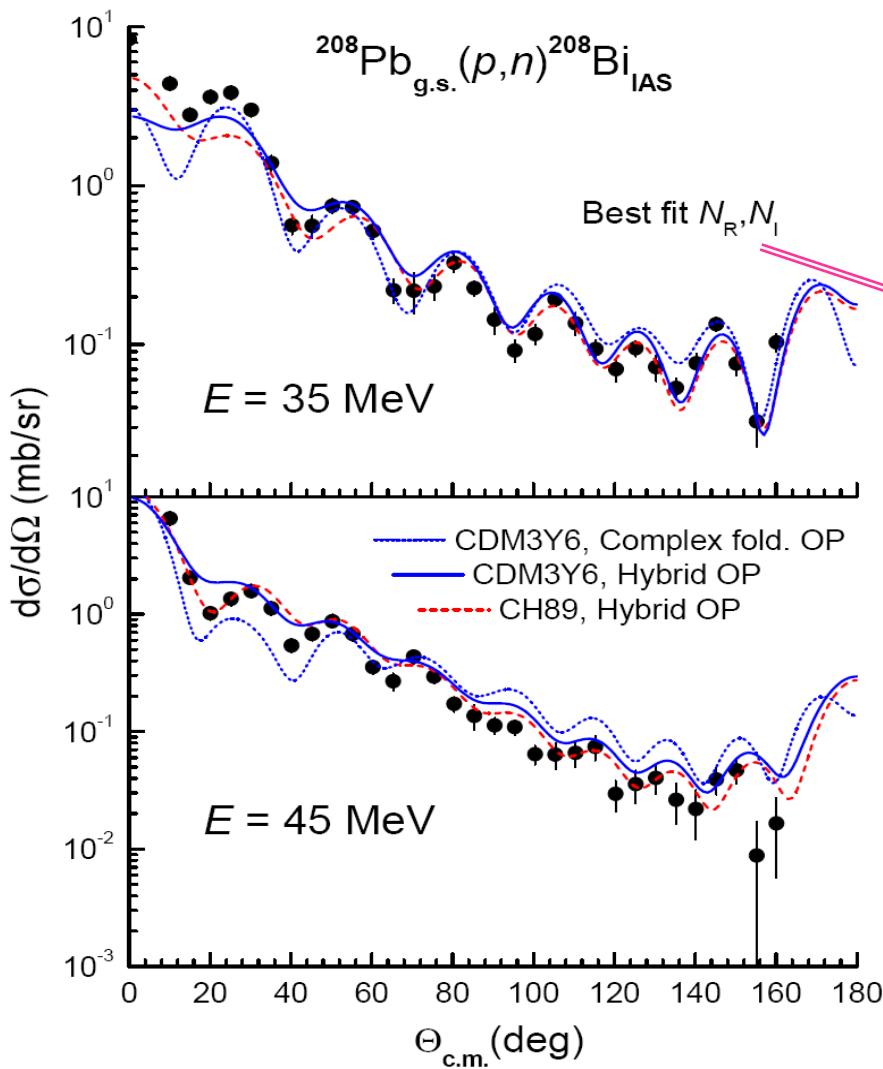
\* Parameters of  $\text{Im}[F_0(\rho, E)]$ ,  $\text{Re}[F_1(\rho, E)]$  and  $\text{Im}[F_1(\rho, E)]$  are adjusted to reproduce the BHF results by J.P. Jeukenne, A. Lejeune and C. Mahaux (**JLM interaction**)  
*Phys. Rev. C 16, 80 (1977).*

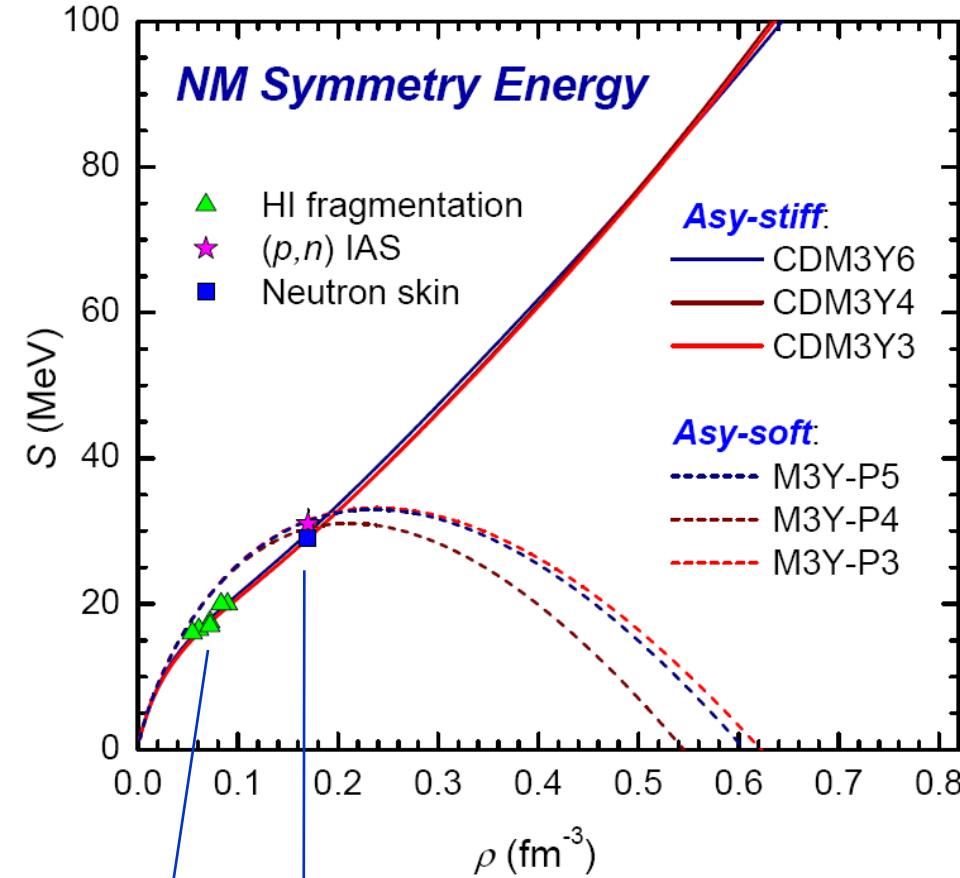
\*  $v_0(r)$  and  $v_1(r)$  are the original M3Y - Paris interaction by N. Anantaraman et al.  
*Nucl. Phys. A398 (1983) 269.*

## STEP II: Adjusting the isospin dependence of the CDM3Y<sub>n</sub> interaction to (p,n) data for IAS excitation !

D.T. Khoa, H.S. Than, and D.C. Cuong, *Phys. Rev. C* **76**, 014603 (2007).

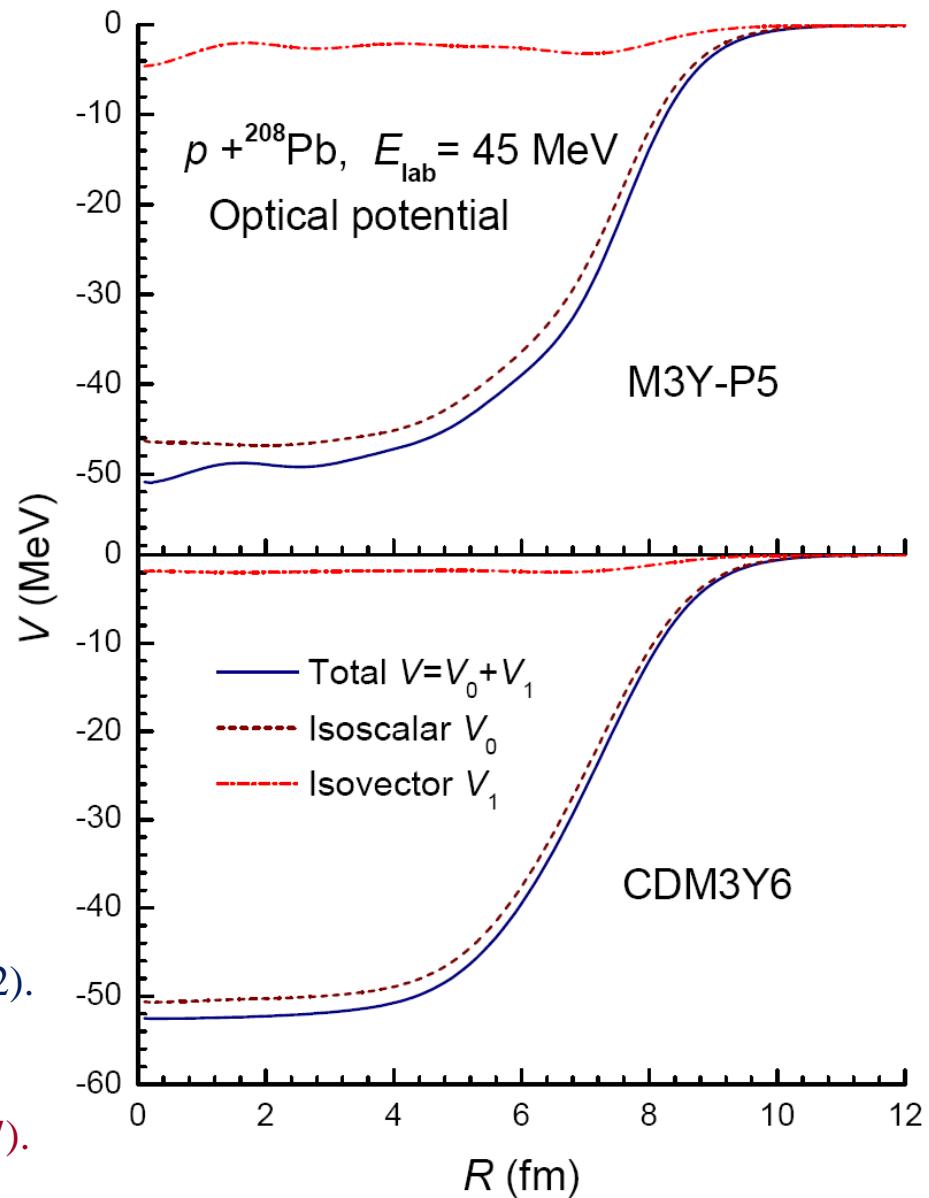
**MSU Data:** R.R. Doering et al. *Phys. Rev. C* **12**, 378 (1975).

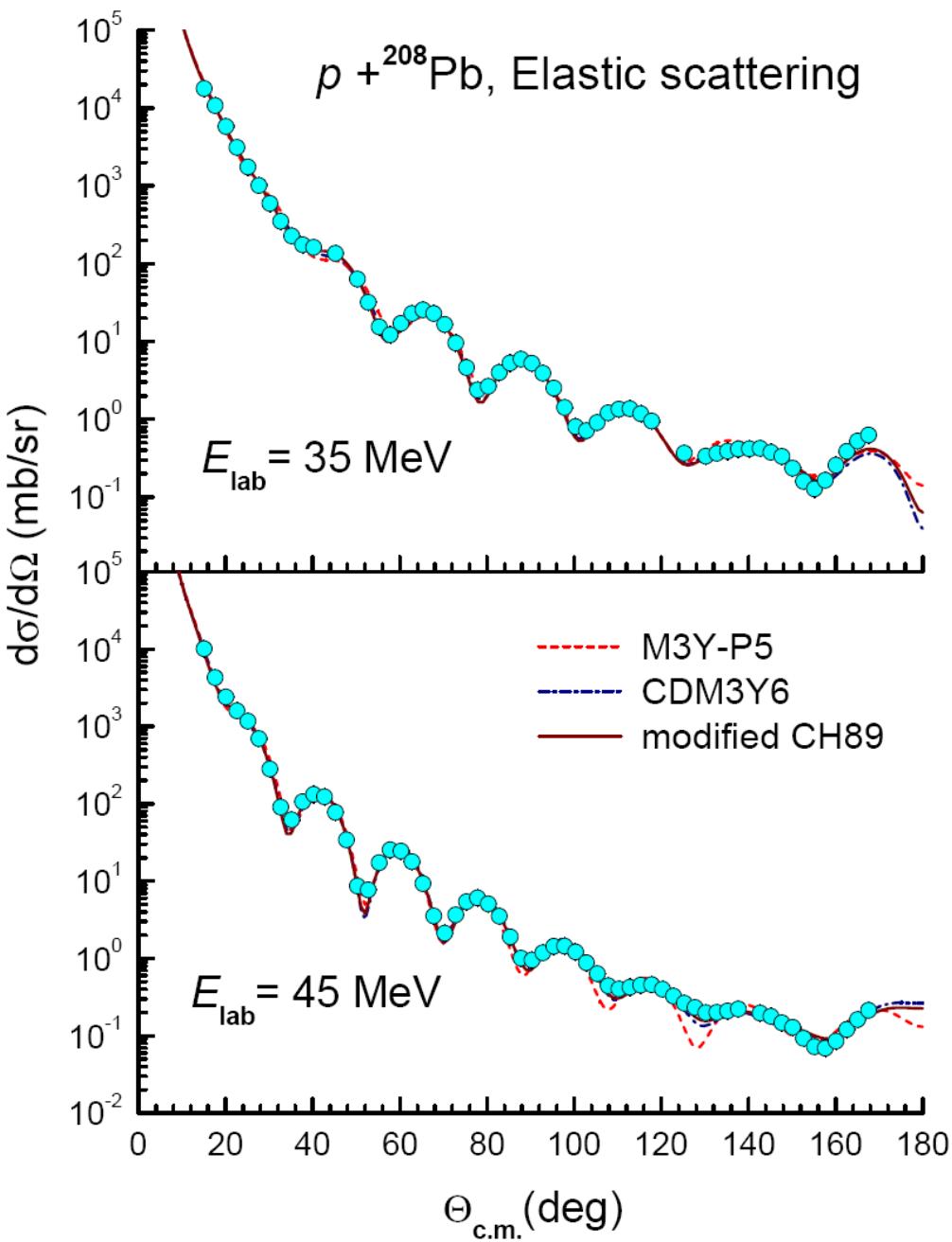




R. J. Furnstahl, *Nucl. Phys. A* **706**, 85 (2002).

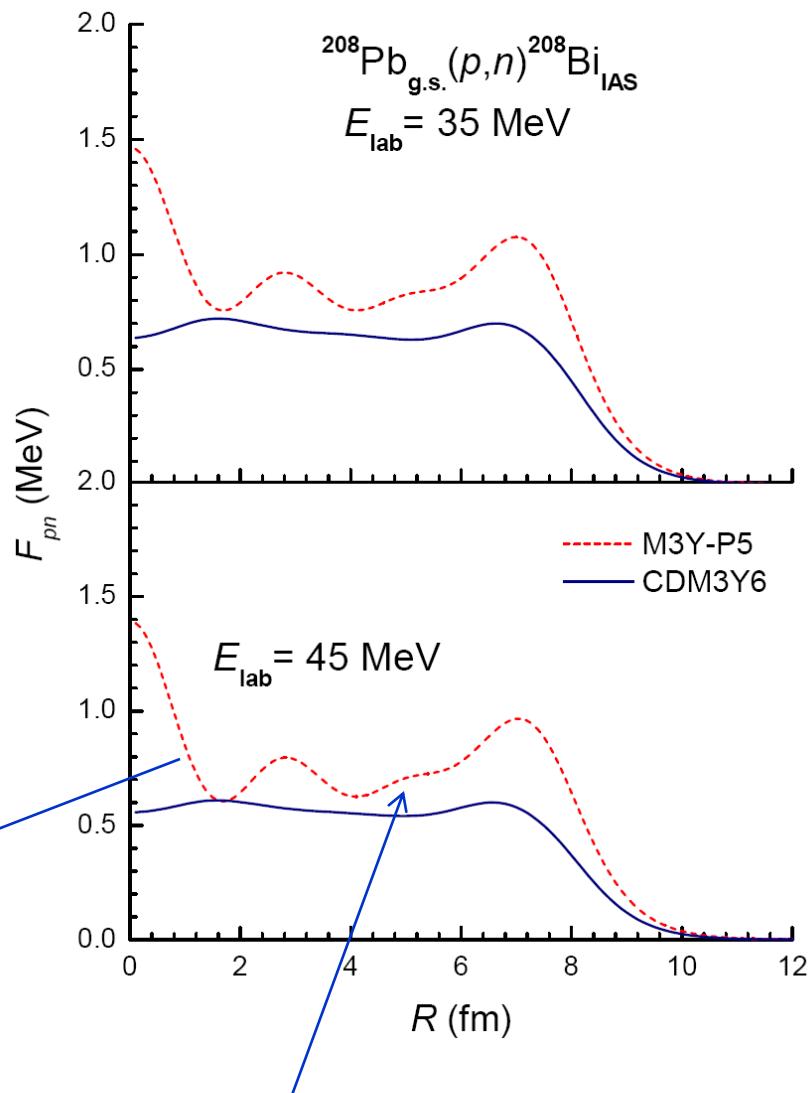
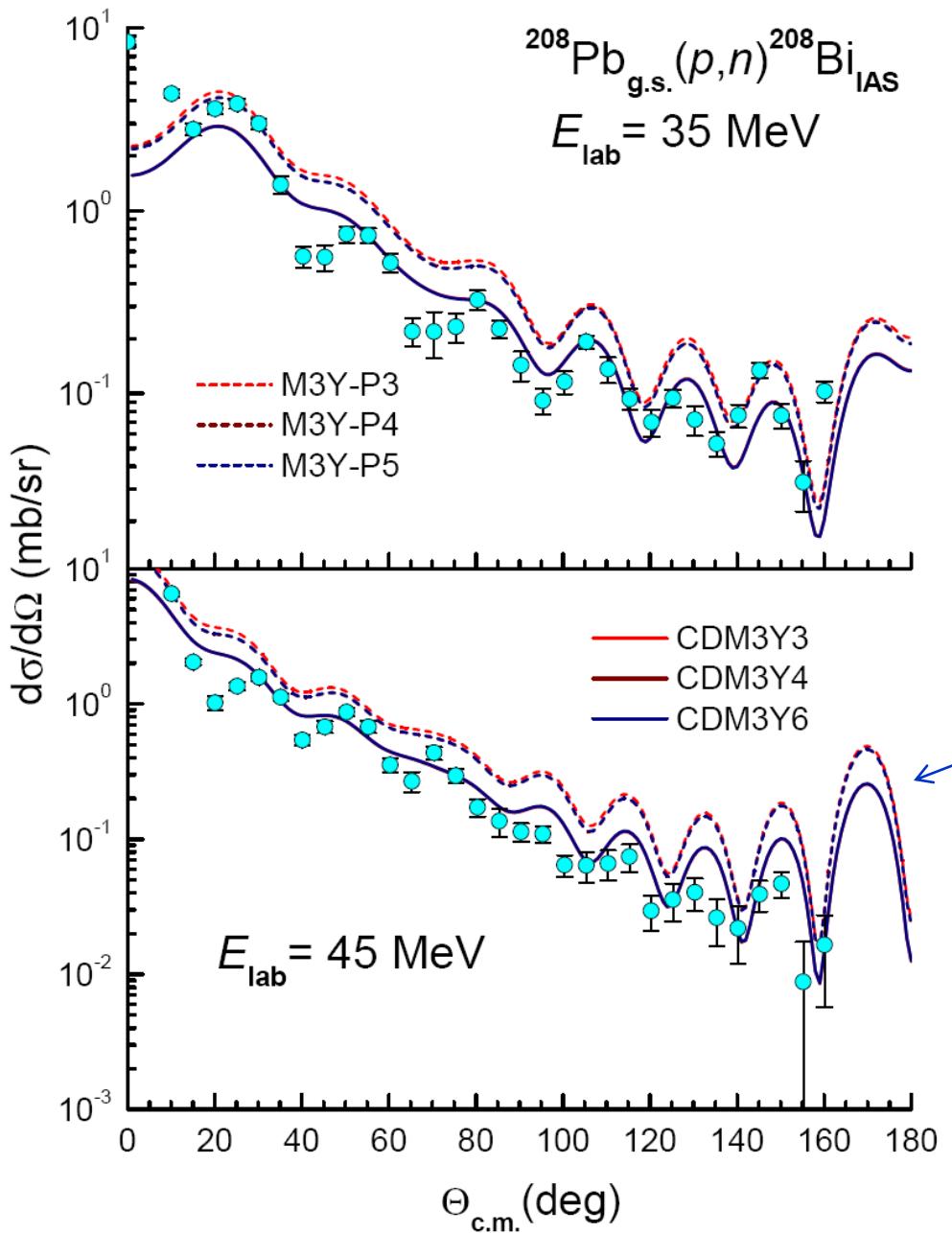
- A. Ono *et al.*, *Phys. Rev. C* **68**, 051601(R) (2003),  
 D. V. Shetty *et al.*, *Phys. Rev. C* **76**, 024606 (2007);  
 D. V. Shetty *et al.*, *NIM Phys. Res. B* **261**, 990 (2007).



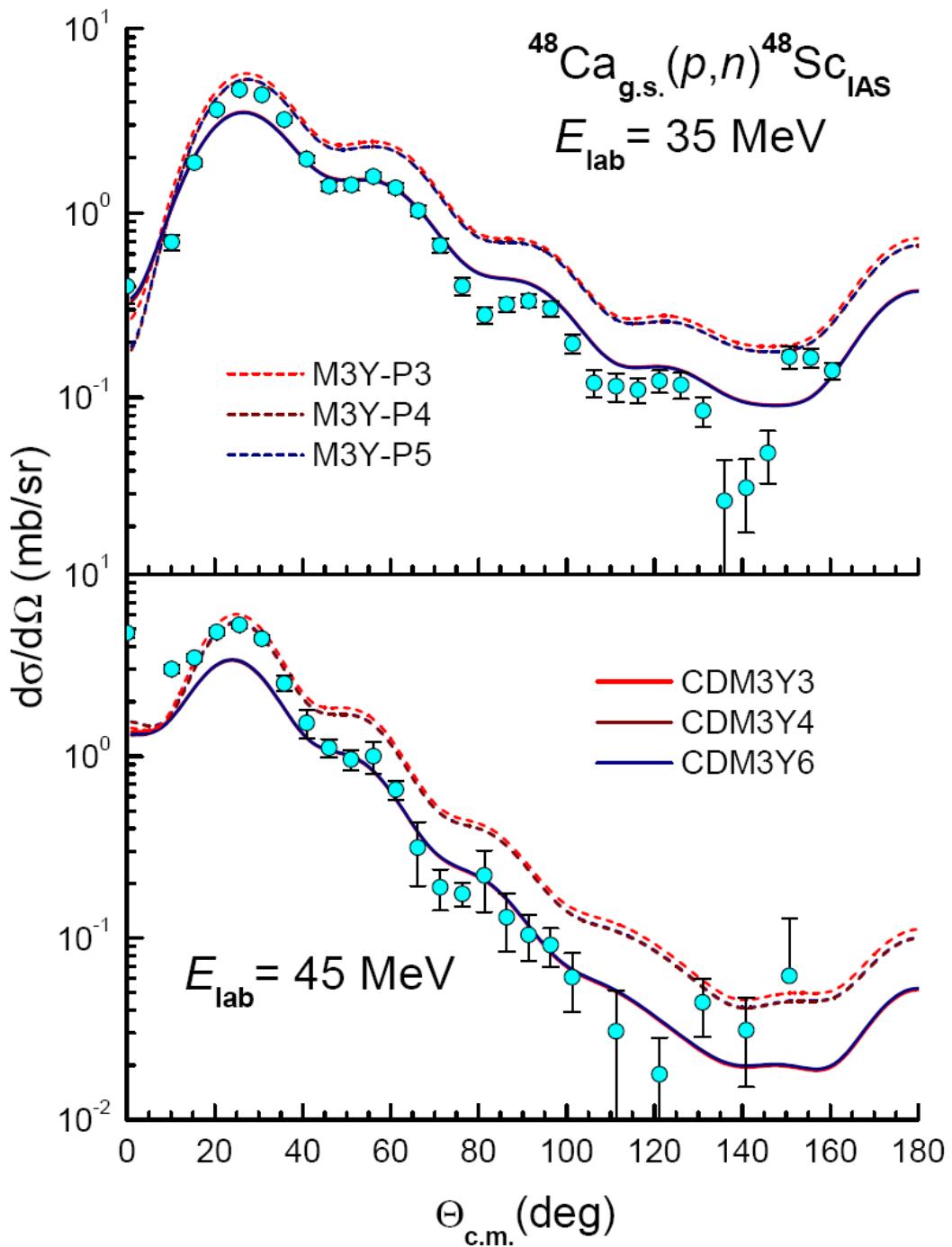


**Data:** W.T.H. van Oers *et al.*,  
*Phys. Rev. C* **10**, 307 (1974)

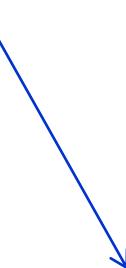
**CH89 OP:** R. L. Varner *et al.*,  
*Phys. Rep.* **201**, 57 (1991)



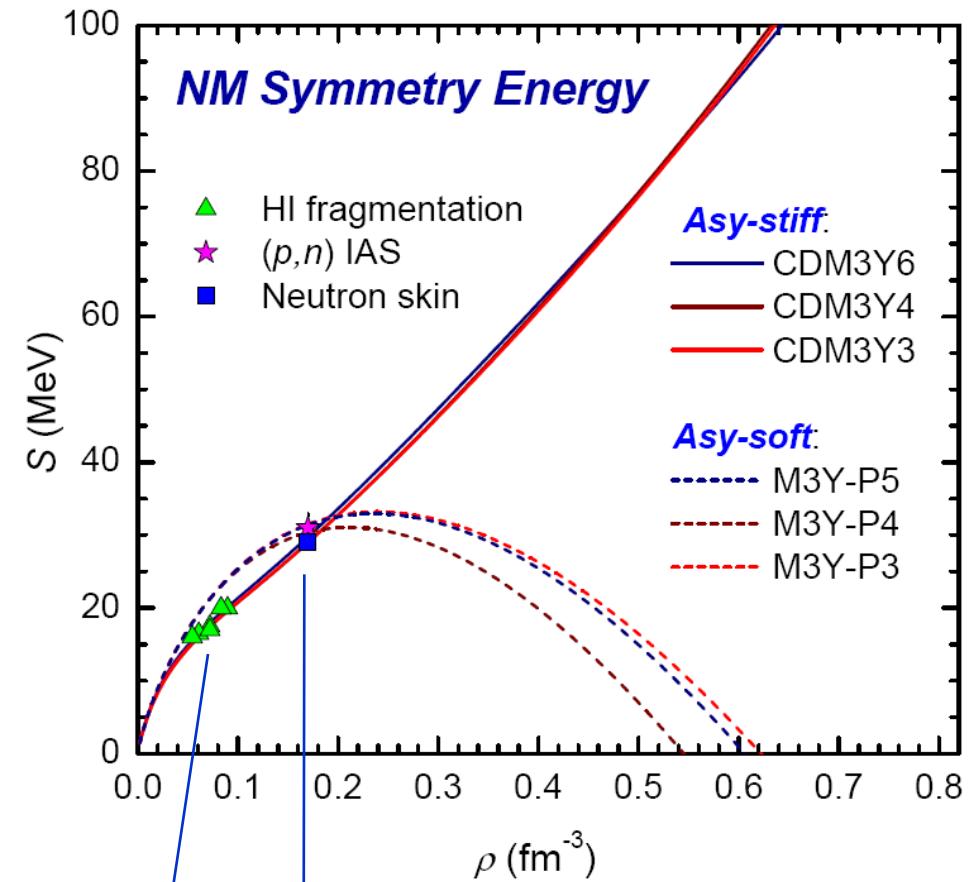
*Difference caused by different  
isospin dependences in CDM3Y6  
and M3Y-P5 interactions !*



**MSU Data:** R.R. Doering *et al.*  
*Phys. Rev. C* **12**, 378 (1975).

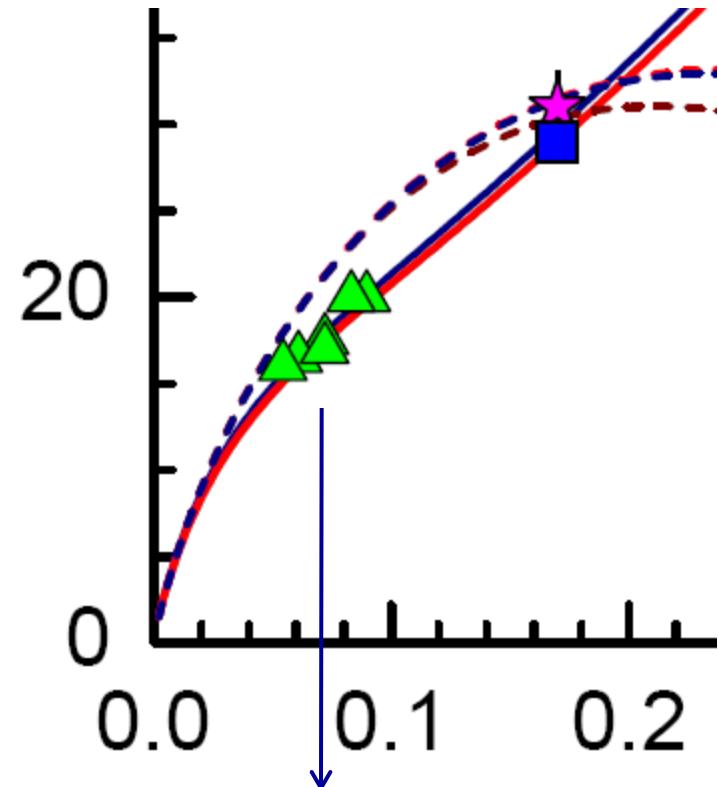


*Stiffness of the Sym. Energy at low barion densities can be probed by  $(p,n)$  IAS data !*



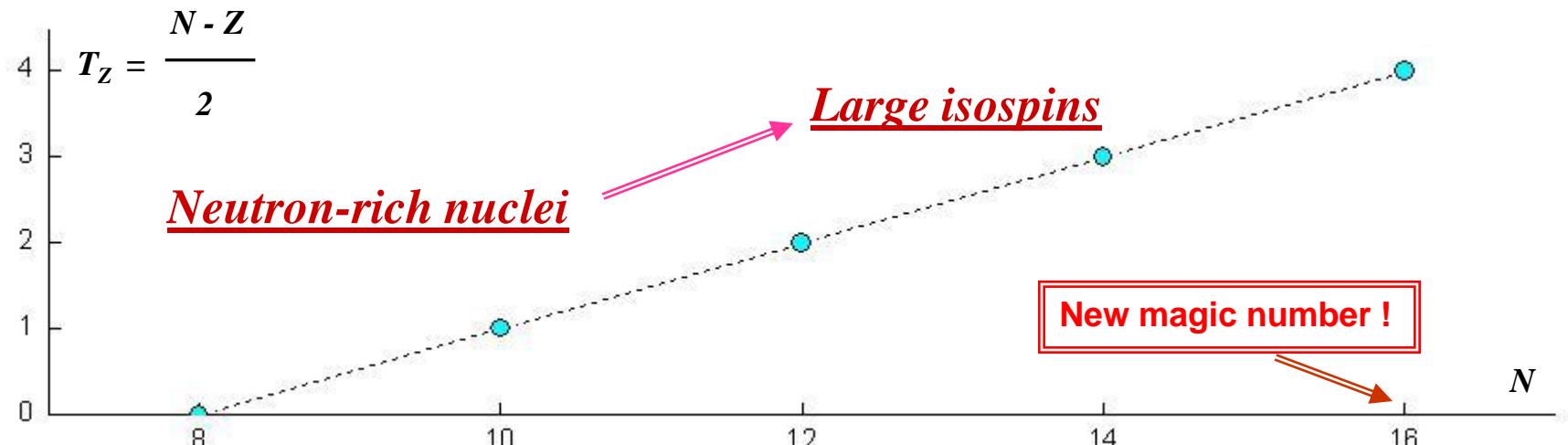
R. J. Furnstahl, *Nucl. Phys. A* **706**, 85 (2002).

- A. Ono *et al.*, *Phys. Rev. C* **68**, 051601(R) (2003),  
 D. V. Shetty *et al.*, *Phys. Rev. C* **76**, 024606 (2007);  
 D. V. Shetty *et al.*, *NIM Phys. Res. B* **261**, 990 (2007).



*Crust - core interface  
 (1<sup>st</sup>-order phase transition from the  
 NS crust to its uniform liquid core)*

Douchin & Haensel,  
*Astronomy & Astrophysics* **380**, 151 (2001).



*Stable  
(99.76%)*



$^{16}\text{O}$

*Stable  
(0.20%)*



$^{18}\text{O}$

$\tau_{1/2} = 13.51\text{ s}$   
(Decay:  $\beta$ )



$^{20}\text{O}$

$\tau_{1/2} = 2.25\text{ s}$   
(Decay:  $\beta$ )



$^{22}\text{O}$

$\tau_{1/2} = 0.061\text{ s}$   
(Decay:  $\beta, n$ )



$^{24}\text{O}$

$\delta = 0$

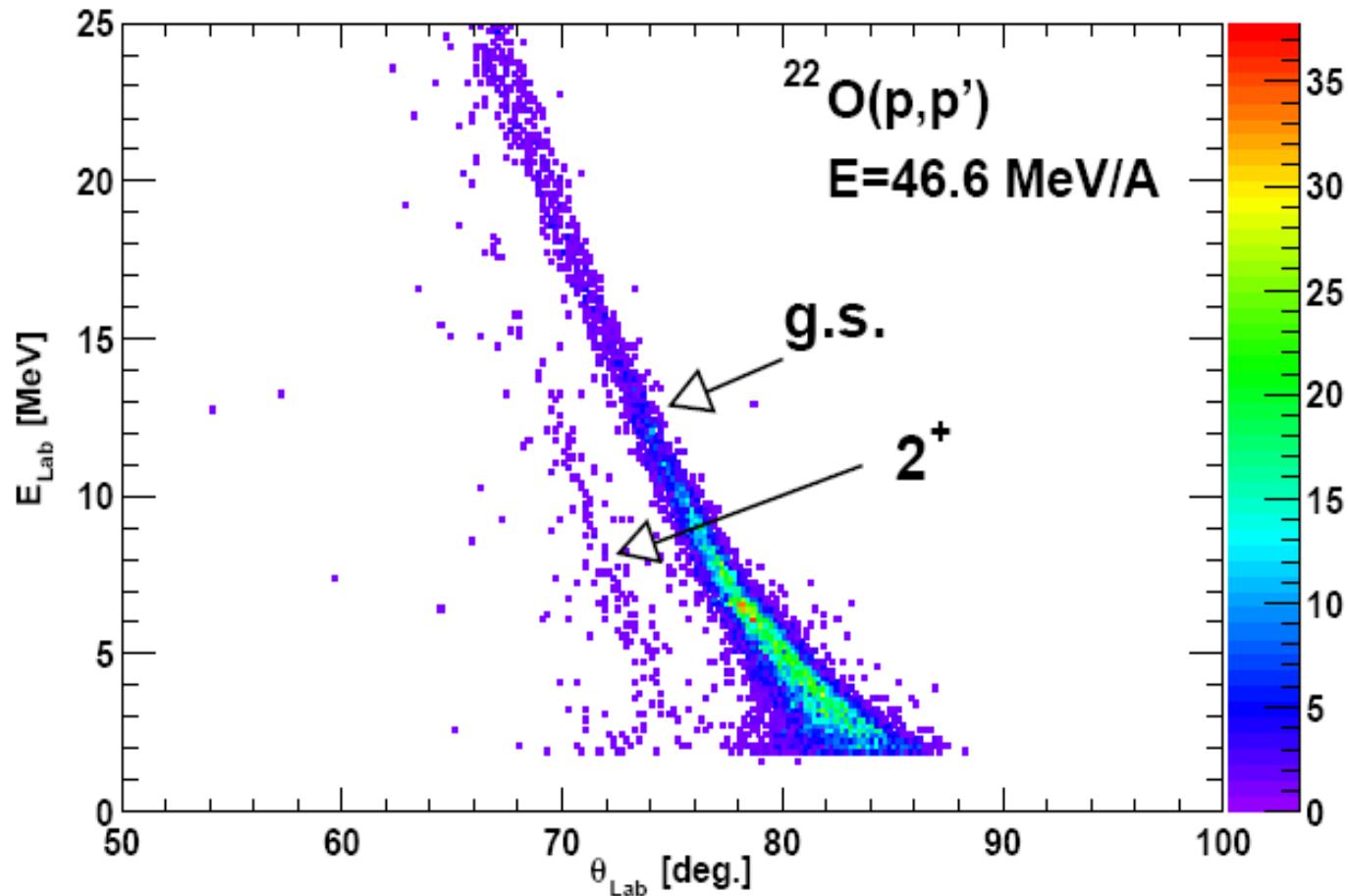
$\delta = 0.11$

$\delta = 0.20$

$\delta = 0.27$

$\delta = 0.33$

$^{36}\text{S}$  primary beam ( $3 \times 10^{12}$  pps) fragmented in the  $^{12}\text{C}$  target located in SISSI device  
⇒ In-flight production of  $^{22}\text{O}$  secondary beam ( $\sim 10^3$  pps) => Hydrogen target.  
⇒ ( $p, p'$ ) events from the detected recoiling protons



## Bohr–Mottelson prescription for the nuclear transition density ( $\lambda \geq 2$ )

$$\rho_\lambda^\tau(r) = -\delta_\tau \frac{d\rho_{g.s.}^\tau(r)}{dr}, \text{ with } \tau = p, n.$$

$\rho_{g.s.}^\tau(r)$  => proton and neutron ground state (g.s.) densities

$\delta_\tau$  => the proton and neutron deformation lengths

### IS and IV transition densities

$$\rho_\lambda^{0(1)}(r) = -\delta_{0(1)} \frac{d[\rho_{g.s.}^n(r) \pm \rho_{g.s.}^p(r)]}{dr}.$$

$\rho_{g.s.}$  calculated in the HFB formalism by Orsay group

M. Grasso, N. Sandulescu, N. Van Giai and R.J. Liotta,  
*Phys. Rev. C* **64**, 064321 (2001).

**IS limit** => neutron ( $\rho_n^p$ ) and proton ( $\rho_p^p$ ) transition densities have the same shape (total g.s. density scaled by  $N/A$  and  $Z/A$ )

$$\Rightarrow \delta_n = \delta_p = \delta_0 = \delta_1$$

$$\frac{M_n}{M_p} = \frac{N}{Z} \quad \text{and} \quad \frac{M_1}{M_0} = \frac{N-Z}{A} = \varepsilon.$$

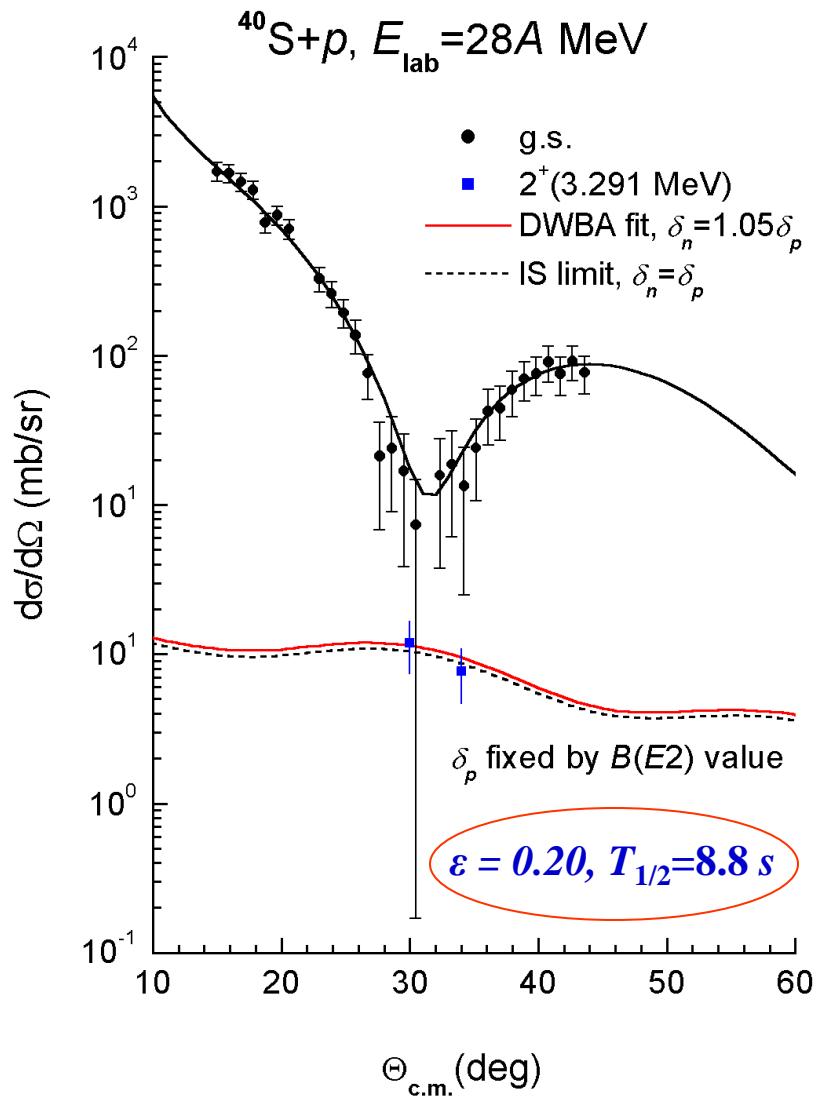
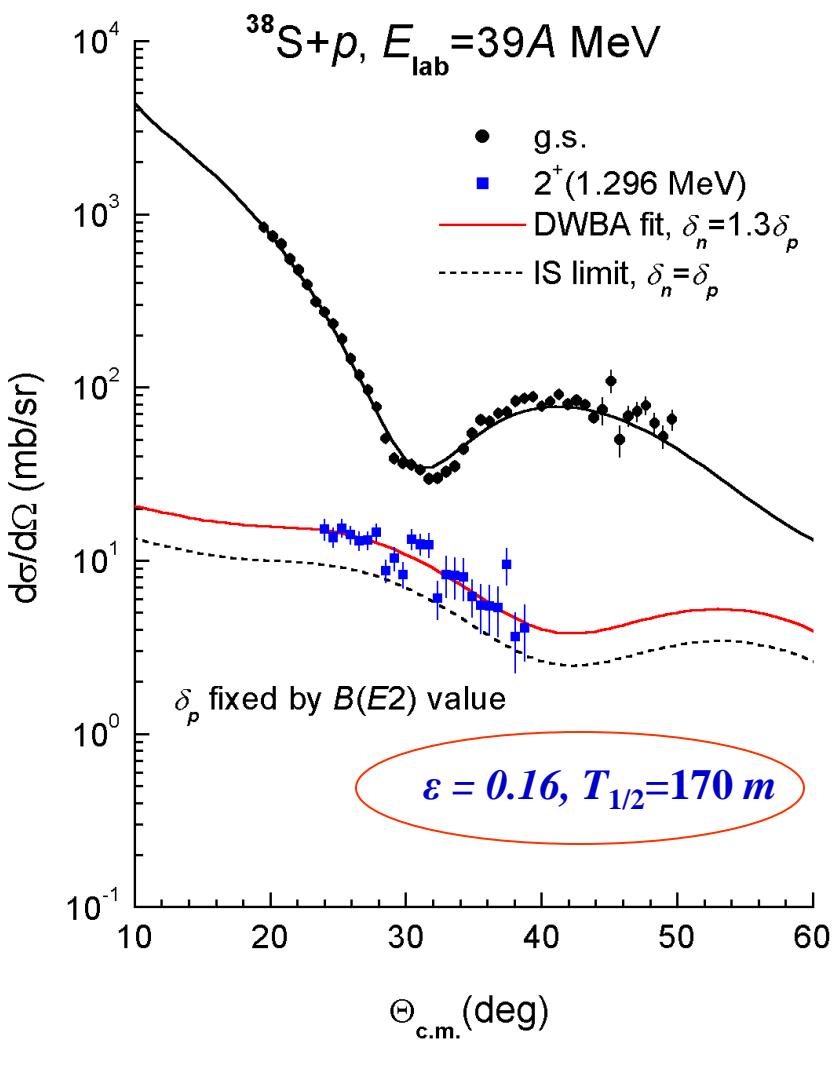
**IV mixing effect** => Difference between  $M_n/M_p$  and  $N/Z$   
(or between  $M_1/M_0$  and  $\varepsilon$ )

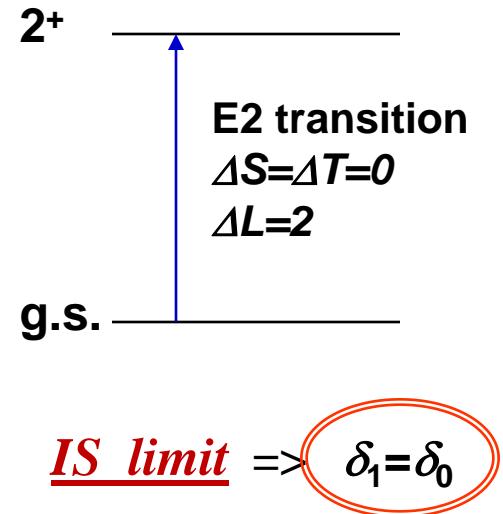
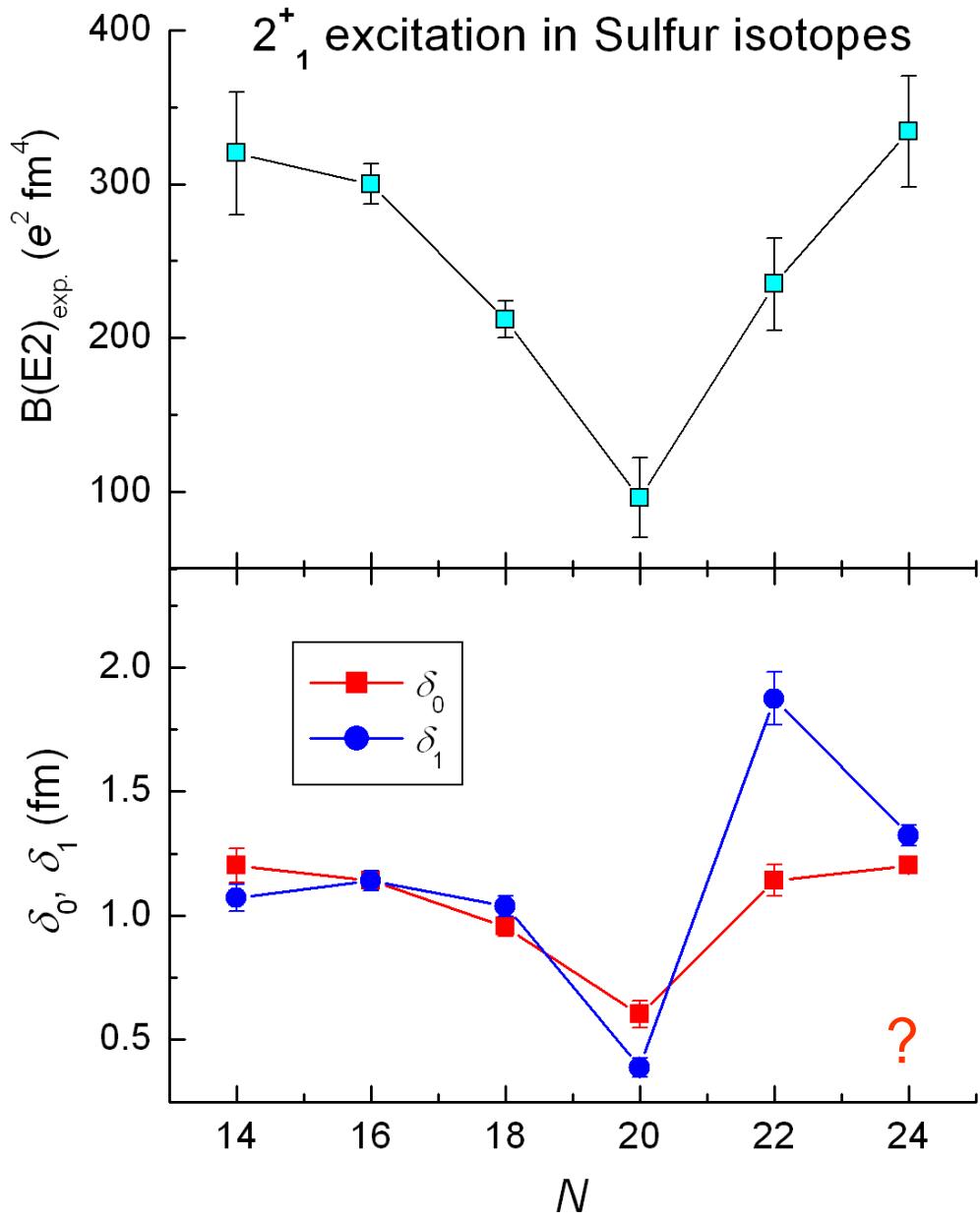
$B(E\lambda\uparrow)_{\text{exp}} = e^2 |M_p|^2 \Rightarrow \delta_p$ ; neutron deformation length  $\delta_n$  is the only parameter determined from DWBA fit to the  $(p,p')$  data.

$\delta_n, \delta_p, M_n, M_p \leq \text{one-to-one correspondence} \Rightarrow \delta_0, \delta_1, M_0, M_1$

**Data:** J. H. Kelley et al., *Phys. Rev. C* **56** (1997) R1206  
and F. Maréchal et al., *Phys Rev. C* **60** (1999) 034615.

**DWBA analysis:** D.T. Khoa, *EPJ Special Topics* **150**, 31 (2007).

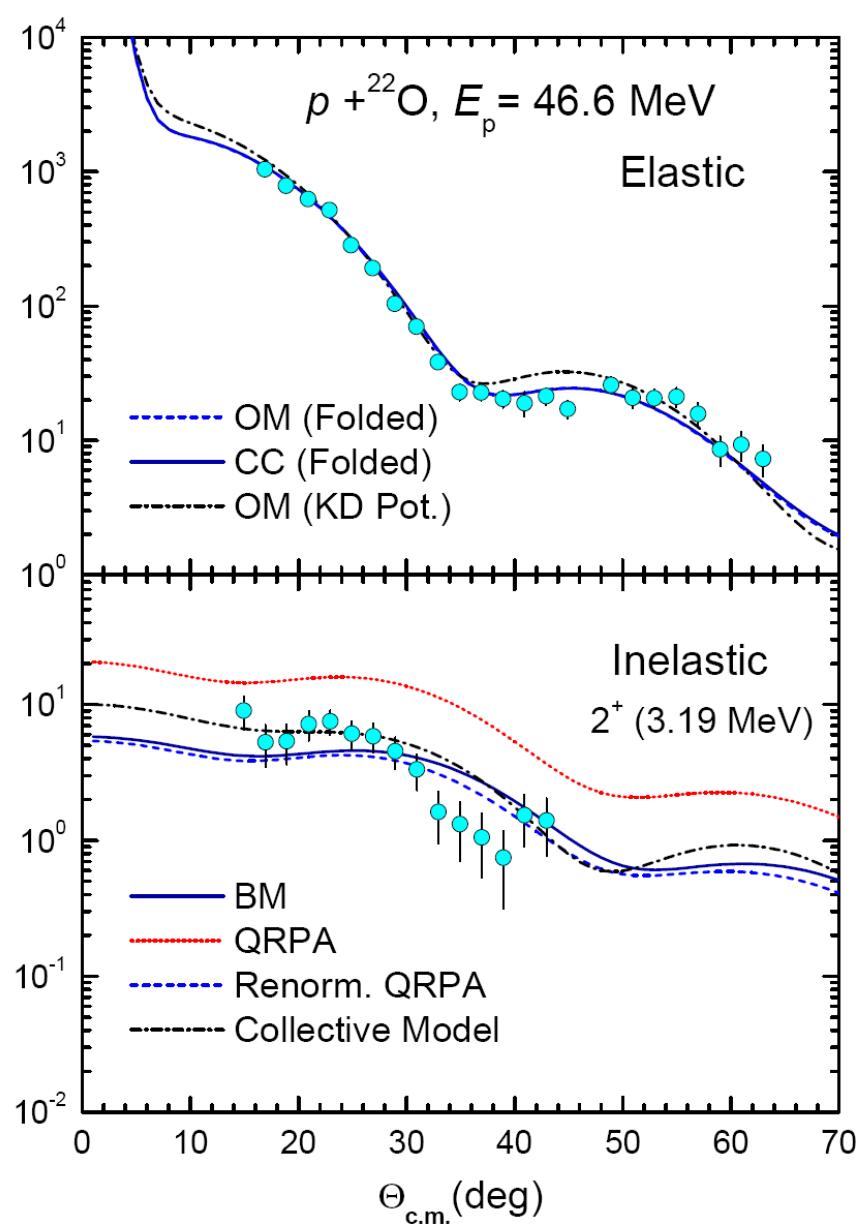
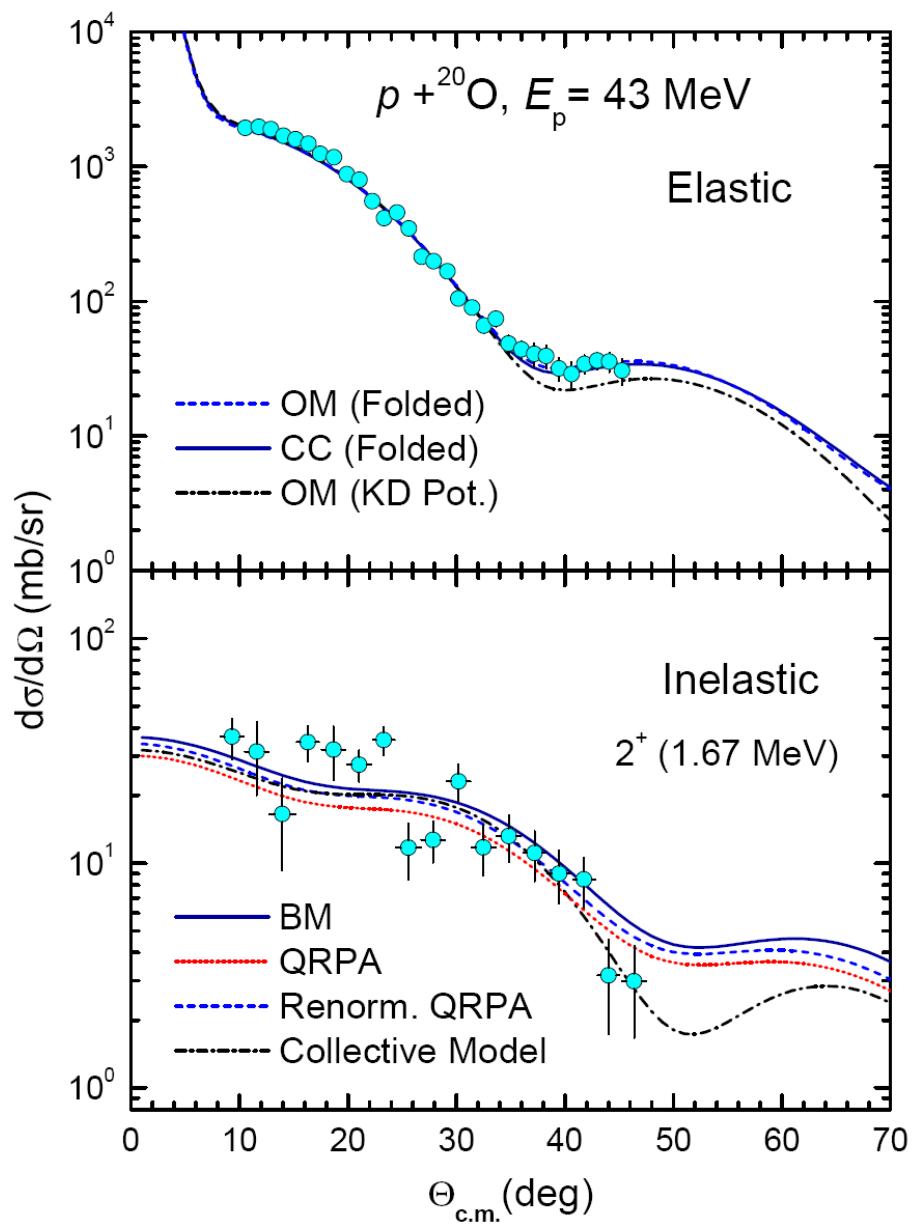


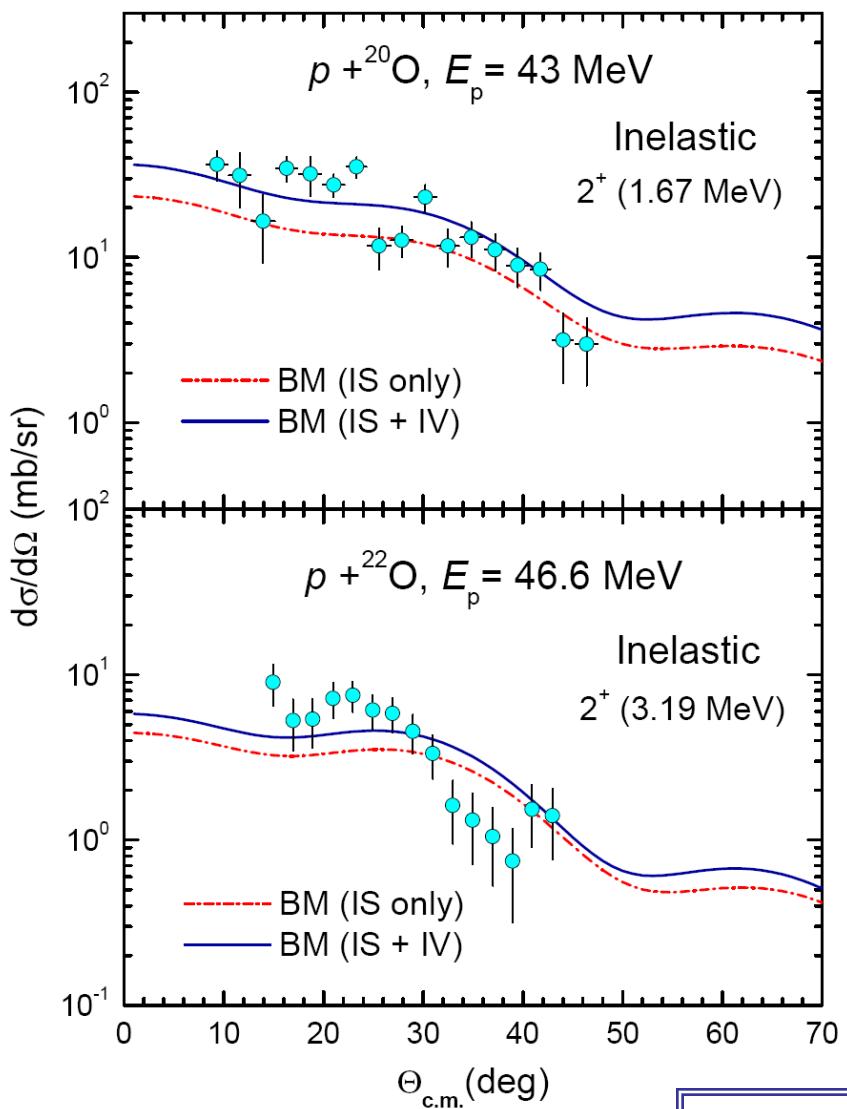


Neutron shell closure  
at  $N = 20$   
 $\Rightarrow$  weak isovector deformation ( $\delta_1 < \delta_0$ )

Data: E. Khan et al., *Phys. Lett. B* **490** (2000) 45; E. Becheva et al., *Phys. Rev. Lett.* **96** (2006) 012501.

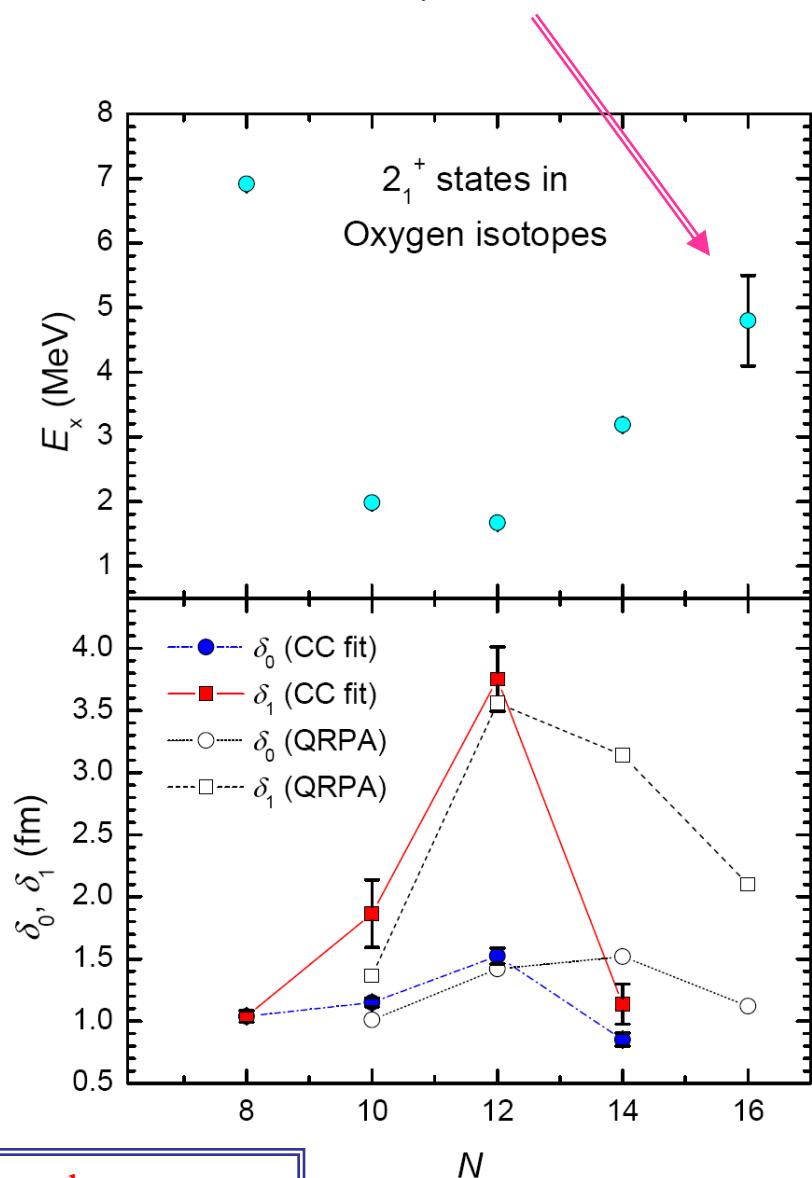
Complex folding + CC analysis: N.D. Chien & D.T. Khoa, *Phys. Rev. C* **79** (2009) 034314.





Weaker IV mixing in  ${}^{22}\text{O}$

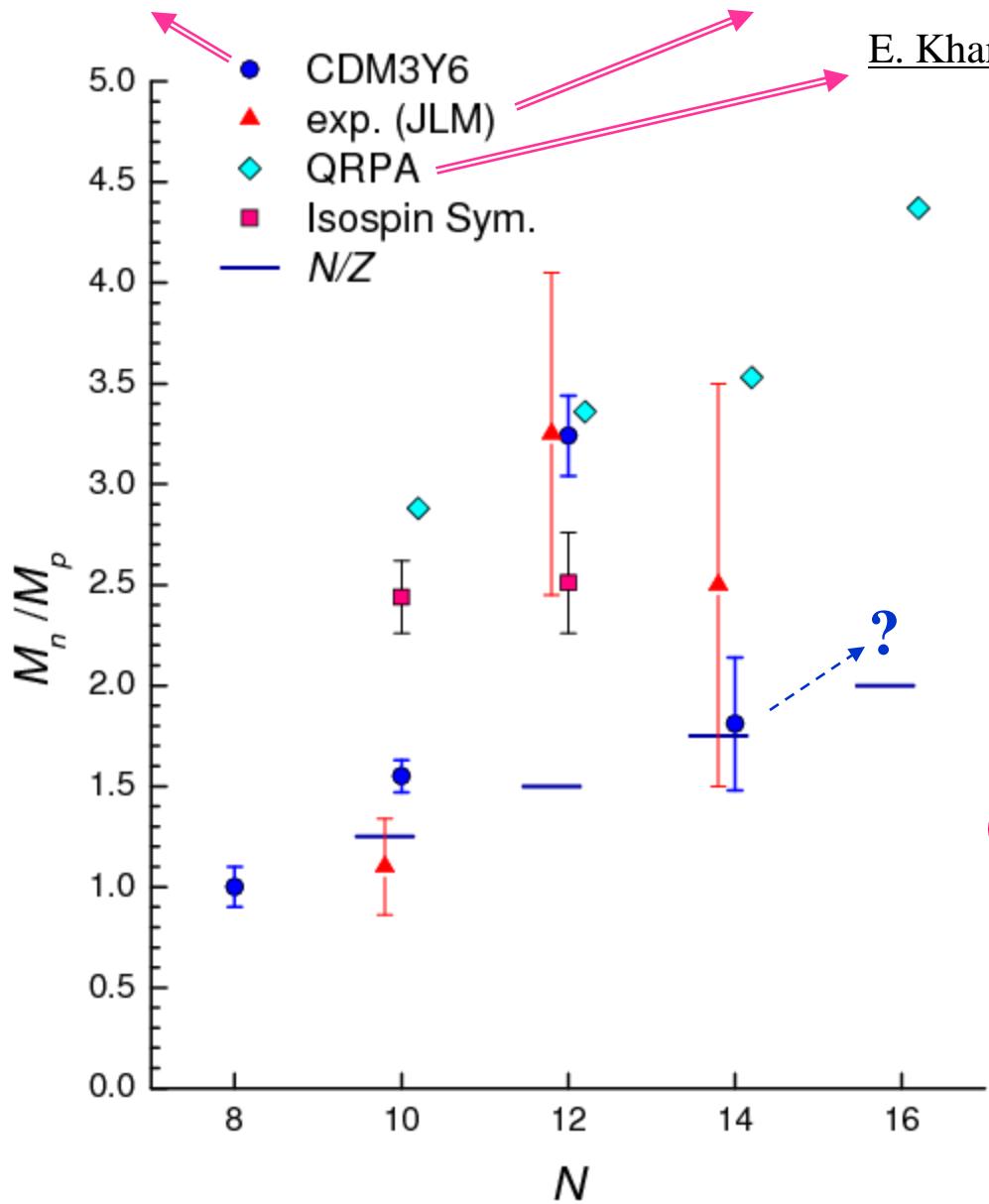
$\sim 4 \text{ MeV gap between}$   
 $2s1/2$  and  $1d3/2$  subshells



New magic number  $N=16$ !

N.D. Chien & D.T. Khoa,  
*Phys. Rev. C* **79** (2009) 034314

E. Khan et al., *Phys. Lett. B* **490** (2000) 45;  
Becheva et al., *Phys. Rev. Lett.* **96** (2006) 012501



E. Khan et al., *Phys. Rev. C* **66** (2002) 024309

*Isospin symmetry*

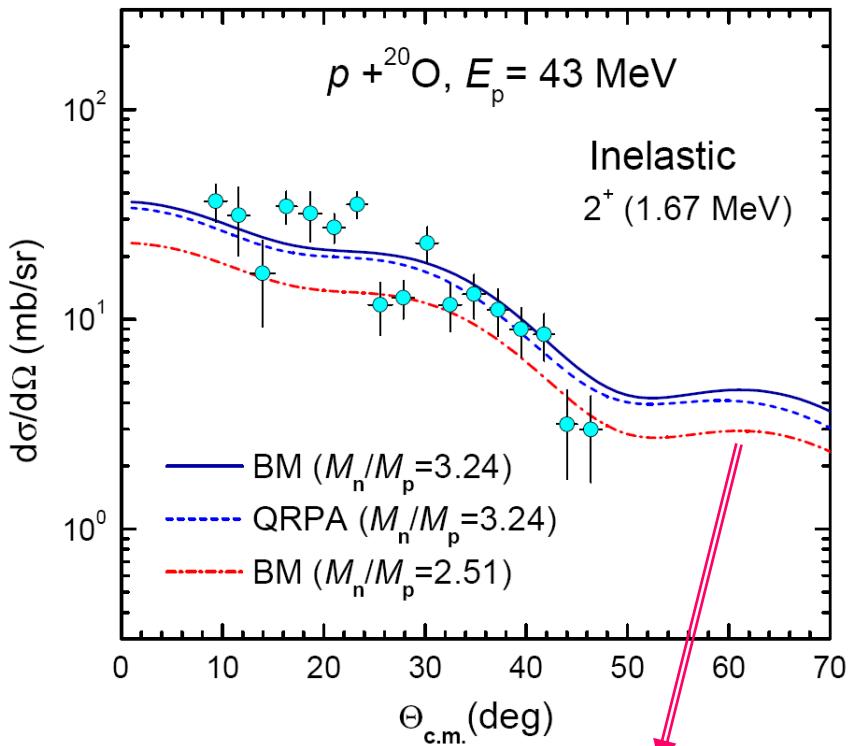
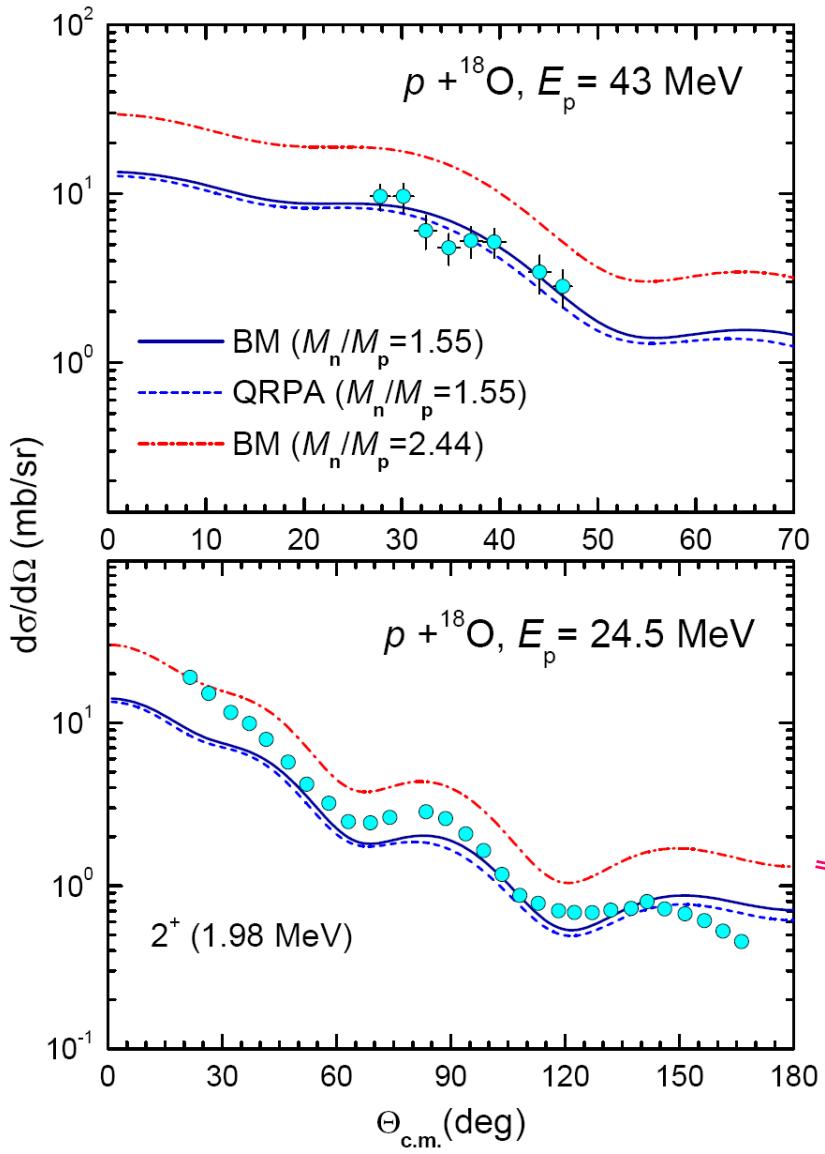
*Charge independence of  $2^+$  excitation  
in members of the T-isospin multiplet*

$$M_p(-T_z) = M_n(T_z)$$

A.M. Bernstein, V.R. Brown, and V.A. Madsen,  
*Phys. Rev. Lett.* **42** (1979) 425.

$B(E2)_{\text{exp}} \Rightarrow M_p$  for  $^{18}\text{Ne}$  and  $^{20}\text{Mg}$   
 $\Rightarrow M_n$  for  $^{18}\text{O}$  and  $^{20}\text{O}$  and vice versa

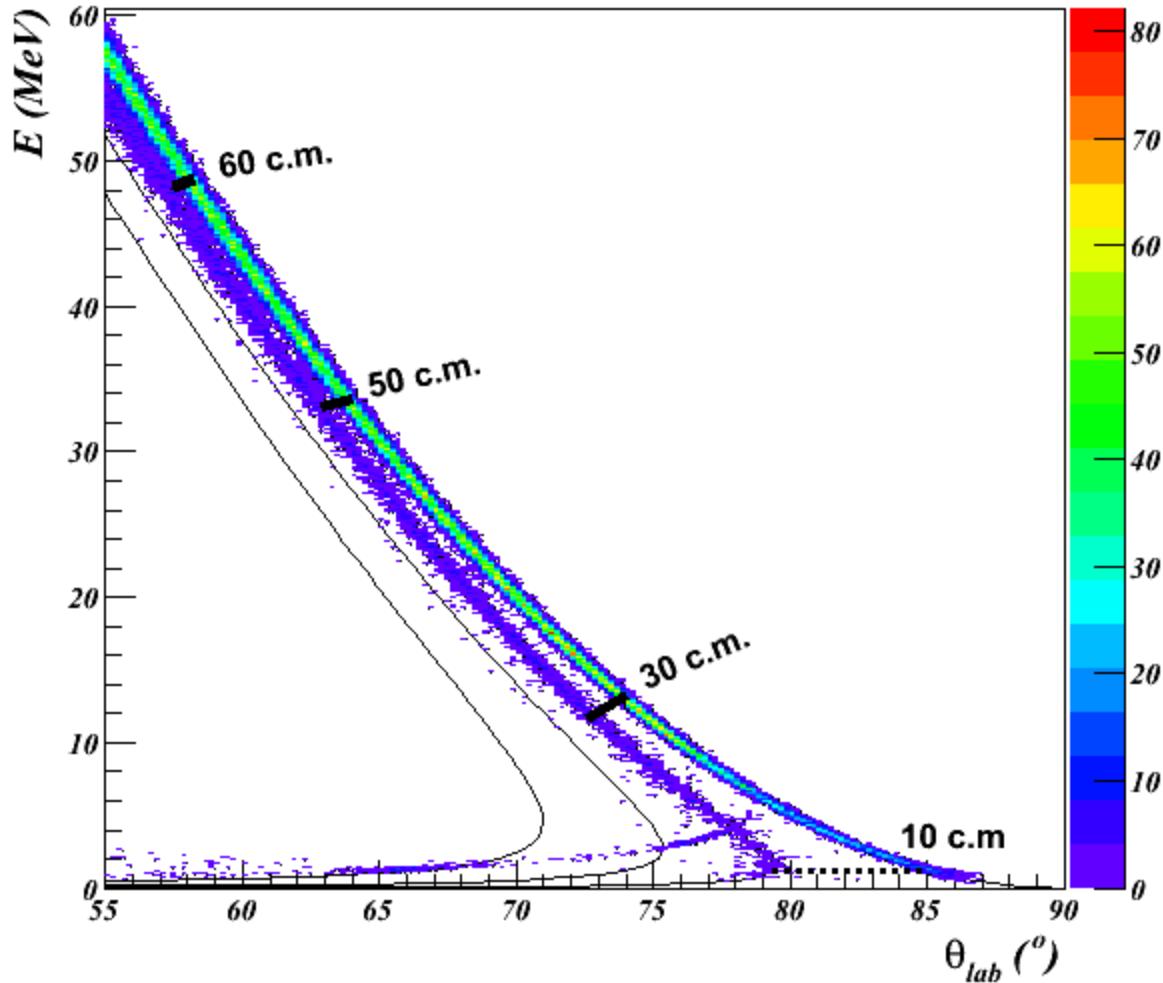
*RIKEN experiment with unstable  $^{20}\text{Mg}$  beam*  
N. Iwasa et al., *Phys. Rev. C* **78** (2008) 024306



*Mirror symmetry in the first  $2^+$  excitation  
 of  $A = 18; T = 1$  and  $A = 20; T = 2$  isobars*

More ( $p,p'$ ) experiments needed !

$^{20}\text{Mg}(p,p')$  at 50 MeV, simulation of recoiled proton spectrum for elastic and inelastic (to  $2^+$  state at 1.6 MeV in  $^{20}\text{Mg}$ ) scattering based on efficiency of MUST2 detector by Valerie Lapoux in September 2009. Not yet measured !!!



*Thank you !*

*Arigato gozaimasu !*

*Cám ơn !*